

Higgs phenomenology beyond the Standard Model

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Physics at the LHC 2011 Perugia, June 6–11

1

Outline

Introduction: how to search for BSM Higgs

Survey of models

Extracting Higgs couplings from LHC data

Conclusions

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Three ways to search for BSM Higgs:

1) Exotic decays of SM-like neutral Higgs

- $H \rightarrow aa \rightarrow 4b$, $2b2\tau$, 4γ

- $H \rightarrow$ dark matter; "hidden valley" states; etc.

2) Production and decay of exotic BSM Higgs-sector states

- H^+ search, e.g. in top decays

- $H^{++} \rightarrow \ell^+ \ell^+$ search

- MSSM A/H via decays to au au

3) Production and decay of BSM Higgs in SM Higgs channels

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3) Production and decay of BSM Higgs in SM Higgs channels

"BSM-ness" manifests via modified coupling(s) to SM particles:

- New particles in the loop

 $gg \rightarrow H$ modified;

 $H \rightarrow \gamma \gamma$ modified

- Sharing of EWSB vev among two or more mass eigenstates VBF \rightarrow H and WH, ZH modified; $H \rightarrow WW, ZZ$ modified $H \rightarrow \gamma\gamma$ can be modified: W in the loop

- Masses of different fermions from different Higgs doublets $gg \rightarrow H$ can be modified; ratios of decays to different fermions modified Can affect all other BRs by changing Higgs total width

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3) Production and decay of BSM Higgs in SM Higgs channels

Examples:

- 4th generation
- MSSM
- more general 2HDMs
- Top-Higgs
- Lee-Wick Standard Model

4th generation

New heavy quarks contribute substantially to $gg \to H$ and $H \to \gamma \gamma$

 $gg \to H$: Loop is independent of m_t for $m_t \gg M_H$. 4th gen t', b' together triple SM amplitude: cross section $9 \times$ SM. (a) (b) $H \rightarrow \gamma \gamma$: н W loop dominates for light \sim Higgs; SM top loop interferes (C) www destructively ($\sim -30\%$). н 4th gen t', b', τ' generally suppress partial width to $\gamma\gamma$.

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4th generation



Black: $gg \rightarrow H \rightarrow WW/SM$ with 4th generation Red: $gg \rightarrow H \rightarrow \gamma\gamma/SM$ with 4th generation Blue: $gg \rightarrow H \rightarrow \gamma\gamma/SM$ with sequential W'Magenta: $gg \rightarrow H \rightarrow \gamma\gamma/SM$ with 4th generation and sequential W'Heather Logan (Carleton U.) Higgs pheno beyond the SM PLHC 2011

7

Squarks, charginos in the loops for ggH and $H\gamma\gamma$

But more important: MSSM has 2 Higgs doublets:

- share in EWSB, mix to form mass eigenstates

$$\bar{g}_{h^0WW} = \bar{g}_{h^0ZZ} = \sin(\beta - \alpha)$$
 $\bar{g}_{H^0WW} = \bar{g}_{H^0ZZ} = \cos(\beta - \alpha)$

- have a nontrivial coupling pattern to fermions

$$-\mathcal{L}_{\mathsf{Yuk}} = Y_u u_R^c H_u Q_L + Y_d d_R^c H_d Q_L + Y_\ell e_R^c H_d L_L + \text{h.c.}$$

$$\bar{g}_{h^{0}t\bar{t}} = \frac{\cos\alpha}{\sin\beta} = \sin(\beta - \alpha) + \cot\beta\cos(\beta - \alpha)$$
$$\bar{g}_{h^{0}b\bar{b}} = \bar{g}_{h^{0}\tau\tau} = -\frac{\sin\alpha}{\cos\beta} = \sin(\beta - \alpha) - \tan\beta\cos(\beta - \alpha)$$

Significant effects when all the Higgs states are relatively light MSSM: $\cos(\beta - \alpha) \simeq \frac{1}{2} \sin 4\beta \frac{m_Z^2}{M_A^2}$

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For a large region of parameter space suppression of the $\gamma\gamma$ mode at the LHC



Suppression still sizable for m_A as large as 500 GeV

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M. Carena, talk at Pheno 2011



LHC reach for the MSSM SM-like Higgs

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3 sigma evidence of the SUSY Higgs responsible for EWSB

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More general 2HDMs

Some similar features to MSSM but generally less constrained

New possibilities for Yukawa structure:

"Lepton-specific" 2HDM: H_q couples to u, d; H_ℓ couples to ℓ can suppress Hqq while enhancing $H\ell\ell$

"Flipped" 2HDM: H_u couples to u and ℓ ; H_d couples to dcan have large Hbb while suppressing $H \to \tau \tau$

More general 2HDMs: Lepton-specific 2HDM benchmark point



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Top-Higgs

Dedicated Higgs doublet (possibly composite) to generate most of top quark mass

Add-on for models of dynamical EWSB: technicolor, 3-site Moose

Top-Higgs doublet has vev $f = v_{SM} \sin \omega$

Top-Higgs particle H_T couples only to $t\bar{t}$, WW, ZZ at tree level

- WW, ZZ couplings suppressed $\sim \sin \omega$
- $t\bar{t}$ coupling enhanced $\sim 1/\sin\omega$

 $gg \rightarrow H_T$ enhanced $\sim 1/\sin^2 \omega$

 H_T BRs significantly modified below WW threshold: no $b\overline{b}$, $\tau\tau$ decays; gg dominates

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Top-Higgs

 $\sin \omega = 0.5 \longrightarrow \sigma(gg \to H_T)/SM \simeq 4$

Tevatron-combined dedicated $gg \rightarrow H \rightarrow WW$ limit from arXiv:1005.3216



Chivukula, Simmons, Coleppa, HEL, & Martin, PRD83, 055013 (2011)

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Lee-Wick Standard Model

Exotic approach to solve hierarchy problem by implementing Pauli-Villars with actual physical fields.

Grinstein, O'Connell, & Wise, PRD 77, 025012 (2008)

Partner fields have opposite sign quadratic Lagrangian terms.

- Seems to violate microscopic causality (!!!)

- Macroscopically all right; unitary at all orders; EW precision constraints satisfied.

Interesting Higgs sector feature is novel mixing structure:

$$\left(\begin{array}{c}h\\\tilde{h}\end{array}\right) = \left(\begin{array}{c}\cosh\theta & \sinh\theta\\\sinh\theta & \cosh\theta\end{array}\right) \left(\begin{array}{c}h_0\\\tilde{h}_0\end{array}\right)$$

Usual $\bar{g}_{hWW}^2 + \bar{g}_{HWW}^2 = 1$ sum rule becomes $\bar{g}_{h_0WW}^2 - \bar{g}_{\tilde{h}_0WW}^2 = 1$. Simultaneously get enhancement of fermion couplings. Only new free parameter is mass ratio $r = M_{h_0}/M_{\tilde{h}_0}$.

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Alvarez, Leskow, & Zurita, arXiv:1104.3496

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To test SM Higgs mechanism, need to measure Higgs couplings.

SM: coupling of Higgs to each SM particle already fixed by known particle masses.

BSM: pattern of deviations from SM expectations characterizes BSM model.



Model-independent Higgs coupling measurements are one of the main selling points of ILC.

Coupling extraction more challenging at LHC due to absence of direct measurement of Higgs production cross section(s).

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Measure event rates at LHC: sensitive to production and decay couplings.

$$\mathsf{Rate}_{ij} = \sigma_i \mathsf{BR}_j = \sigma_i \frac{\mathsf{\Gamma}_j}{\mathsf{\Gamma}_{\mathsf{tot}}}$$

Main difficulty: "flat direction" in the fit.

Allow an unobserved decay mode while simultaneously increasing all couplings by a factor a:

$$\text{Rate}_{ij} = a^2 \sigma_i^{\text{SM}} \frac{a^2 \Gamma_j^{\text{SM}}}{a^2 \Gamma_{\text{tot}}^{\text{SM}} + \Gamma_{\text{new}}}$$

Ways to deal with this:

- assume no unobserved decays
- assume HWW, HZZ couplings no larger than in SM
- include direct measurement of Higgs width (heavier masses)*

*new

Get ratios of Higgs couplings-squared from taking ratios of rates. Full coupling extraction: assume no unexpected decay channels, assume $b\bar{b}/\tau\tau = SM$ value. $M_H = 100-190$ GeV

Zeppenfeld, Kinnunen, Nikitenko, Richter-Was, PRD62, 013009 (2000); Les Houches 1999

Add $t\bar{t}H$, $H \rightarrow \tau\tau$ channel to improve $t\bar{t}H$ constraint. $M_H = 110-180 \text{ GeV}$ Belyaev & Reina, JHEP0208, 041 (2002)

Fit assuming WWH, ZZH couplings bounded from above by SM value. $M_H = 110-190$ GeV Dührssen, Heinemeyer, HEL, Rainwater, Weiglein, & Zeppenfeld, PRD70, 113009 (2004)

More careful analysis of probability density and correlations, using updated expt studies. Assume no unexpected decay channels. $M_H = 120$ GeV Lafaye, Plehn, Rauch, D. Zerwas, & Dührssen, JHEP0908, 009 (2009)

New approach for heavier Higgs ($\gtrsim 190 \text{ GeV}$) using direct Higgs width measurement from $H \rightarrow ZZ \rightarrow 4\ell$ lineshape. $M_H = 190 \text{ GeV}$ HEL & Salvail, in preparation

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Dührssen, Heinemeyer, HEL, Rainwater, Weiglein, & Zeppenfeld, PRD70, 113009 (2004)

- 10%–50%+ uncertainties on couplings-squared.
- Systematic & theory uncertainties are important.

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Lafaye, Plehn, Rauch, D. Zerwas, & Dührssen, JHEP 0908, 009 (2009)

- Correlations in parameters are also important
- Results depend on what coupling freedom is assumed (here: no additional Hgg, $H\gamma\gamma$ contributions)

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Can we make model-independent measurements?

$$\text{Rate}_{ij} = a^2 \sigma_i^{\text{SM}} \frac{a^2 \Gamma_j^{\text{SM}}}{a^2 \Gamma_{\text{tot}}^{\text{SM}} + \Gamma_{\text{new}}}$$

Consider extraction of Higgs couplings when Higgs total width is a directly measurable observable.



CMS TDR (2006), Vol. 2 (Physics), chap. 10

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Total width: 17.6% for 30 fb⁻¹, 9.6% for 100 fb⁻¹ CMS TDR (2006), Vol. 2 (Physics)

Rates:

Production	Decay	$30~{ m fb}^{-1}$	$100~{ m fb}^{-1}$	"contamination"	
GF	$ZZ ightarrow 4\ell$	14%	7.9%	${\sf VBF}\sim 14\%$	a
VBF	$ZZ ightarrow 4\ell$	24%	13%	${\sf GF}\sim 21\%$	a
GF	$WW ightarrow \ell \ell p_T^{miss}$	9.6%	5.3%	VBF $\sim 2.8\%$	a
VBF	$WW ightarrow e \mu p_T^{miss}$	14%	7.6%	${\sf GF}\sim 7.8\%$	a
VBF	$WW ightarrow (ee, \mu \mu) p_T^{miss}$	15%	8.1%	${ m GF}\sim 7.2\%$	a
VBF	$WW ightarrow \ell u j j$	16%	8.9%	(none)	b
	1				

^a Dührssen, ATL-PHYS-2003-030 ^b Pi et al, CMS-NOTE-2006-092

- All uncertainties statistical only.

- All studies 30 fb⁻¹ at 1 detector at 14 TeV; we scale by $\sqrt{3/10}$ for 100 fb⁻¹ estimate.

Parametrization of new physics:

 $\Gamma_{\text{tot}} = \Gamma_W + \Gamma_Z + \Gamma_{\text{new}}$ $\Gamma_W = \bar{g}_W^2 \Gamma_W^{\text{SM}}$ $\Gamma_Z = \bar{g}_Z^2 \Gamma_Z^{\text{SM}}$ $\sigma_{\rm GF} = \overline{g}_q^2 \sigma_{\rm GF}^{\rm SM} \qquad \sigma_{\rm VBF} = [0.73 \overline{g}_W^2 + (1 - 0.73) \overline{g}_Z^2] \sigma_{\rm VBF}^{\rm SM}$ Heather Logan (Carleton U.) Higgs pheno beyond the SM PLHC 2011



HEL & Salvail, in preparation

$$\delta \overline{g}_V^2 \simeq 8\%, \ \delta \overline{g}_g^2 \simeq 11\%, \ \Gamma_{\rm new}/\Gamma_{\rm tot}^{\rm SM} \lesssim 24\% \ {
m at} \ 95\% \ {
m CL}$$

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HEL & Salvail, in preparation

$$\begin{split} &\delta \bar{g}_V^2 \simeq 4.5\%, \ \delta \bar{g}_g^2 \simeq 5.8\%, \ \Gamma_{\text{new}}/\Gamma_{\text{tot}} \lesssim 12\% \text{ at } 95\% \text{ CL} \\ & \text{Heather Logan (Carleton U.)} & \text{Higgs pheno beyond the SM} & \text{PLHC 2011} \end{split}$$

Conclusions

SM Higgs discovery or exclusion is imminent!

A signal in one of the SM Higgs search channels will have immediate impact on BSM Higgs scenarios.

Higgs coupling measurements are key to understanding structure of Higgs sector.