

# $H \rightarrow hh$ in the Georgi-Machacek model

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## Introduction: motivating the Georgi-Machacek model

Can we constrain the possibility that “exotic” Higgs fields (isospin  $> 1/2$ ) contribute to electroweak symmetry breaking?

Generically this is very strongly constrained by the  $\rho$  parameter:

$$\rho \equiv \frac{\text{weak neutral current}}{\text{weak charged current}} = \frac{(g^2 + g'^2)/M_Z^2}{g^2/M_W^2} = \frac{v_\phi^2 + a\langle X^0 \rangle^2}{v_\phi^2 + b\langle X^0 \rangle^2}$$

$$a = 4 [T(T + 1) - Y^2] c$$

$$b = 8Y^2$$

$$Q = T^3 + Y; \text{ SM doublet: } Y = 1/2$$

Expt:  $\rho = 1.00039 \pm 0.00019$  (2018 PDG)

Need to do some model-building; otherwise  $v_{\text{exotic}} \ll v_{\text{doublet}}$ .

There are only two known approaches:

1) Use the septet  $(T, Y) = (3, 2)$ :  $\rho = 1$  by accident!

Doublet  $(\frac{1}{2}, \frac{1}{2}) +$  septet  $(3, 2)$ : **Scalar septet model**

Hisano & Tsumura, 1301.6455; Kanemura, Kikuchi & Yagyu, 1301.7303

2) Use global  $SU(2)_L \times SU(2)_R$  imposed on the scalar potential

Global  $SU(2)_L \times SU(2)_R \rightarrow$  custodial  $SU(2)$  ensures tree-level  $\rho = 1$

Doublet + triplets  $(1, 0) + (1, 1)$ : **Georgi-Machacek model**

Georgi & Machacek 1985; Chanowitz & Golden 1985

Doublet + quartets  $(\frac{3}{2}, \frac{1}{2}) + (\frac{3}{2}, \frac{3}{2})$ : **Generalized Georgi-**

Doublet + quintets  $(2, 0) + (2, 1) + (2, 2)$ : **Machacek models**

Doublet + sextets  $(\frac{5}{2}, \frac{1}{2}) + (\frac{5}{2}, \frac{3}{2}) + (\frac{5}{2}, \frac{5}{2})$ :

Galison 1984; Robinett 1985; HEL 1999; Chang et al 2012; HEL & Rentala 2015

Larger than sextets  $\rightarrow$  too many large multiplets, violates perturbativity

Can also have duplications, combinations  $\rightarrow$  ignore that here.

Georgi-Machacek model Georgi & Machacek 1985; Chanowitz & Golden 1985

SM Higgs (bi-)doublet + triplets (1, 0) + (1, 1) in a bi-triplet:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix} \quad X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{++*} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

Global  $SU(2)_L \times SU(2)_R \rightarrow$  custodial symmetry  $\langle \chi^0 \rangle = \langle \xi^0 \rangle \equiv v_\chi$

Physical spectrum:

Bi-doublet:  $2 \otimes 2 \rightarrow 1 \oplus 3$

Bi-triplet:  $3 \otimes 3 \rightarrow 1 \oplus 3 \oplus 5$

- Two custodial singlets mix  $\rightarrow h, H$   $m_h, m_H, \text{angle } \alpha$   
Usually identify  $h = h(125)$
- Two custodial triplets mix  $\rightarrow (H_3^+, H_3^0, H_3^-)$   $m_3$  + Goldstones  
Phenomenology very similar to  $H^\pm, A^0$  in 2HDM Type I,  $\tan \beta \rightarrow \cot \theta_H$
- Custodial fiveplet  $(H_5^{++}, H_5^+, H_5^0, H_5^-, H_5^{--})$   $m_5$   
Fermiophobic;  $H_5 VV$  couplings  $\propto s_H \equiv \sqrt{8}v_\chi/v_{SM}$   
 $s_H^2 \equiv$  exotic fraction of  $M_W^2, M_Z^2$

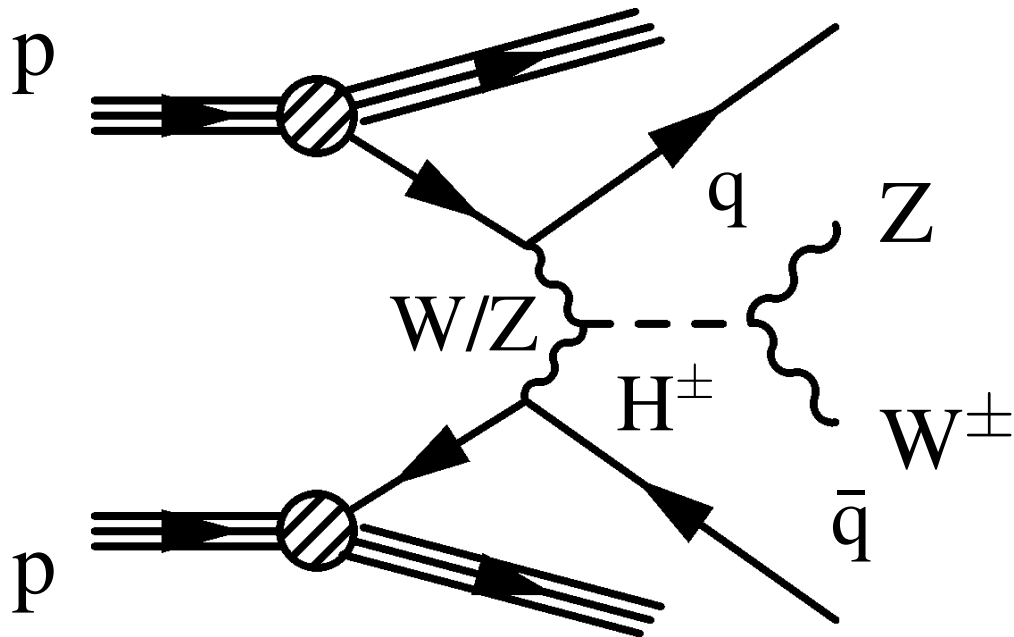
Usual LHC searches:

$$\text{VBF} \rightarrow H_5^{\pm\pm} \rightarrow W^\pm W^\pm$$

VBF + like-sign dileptons + MET

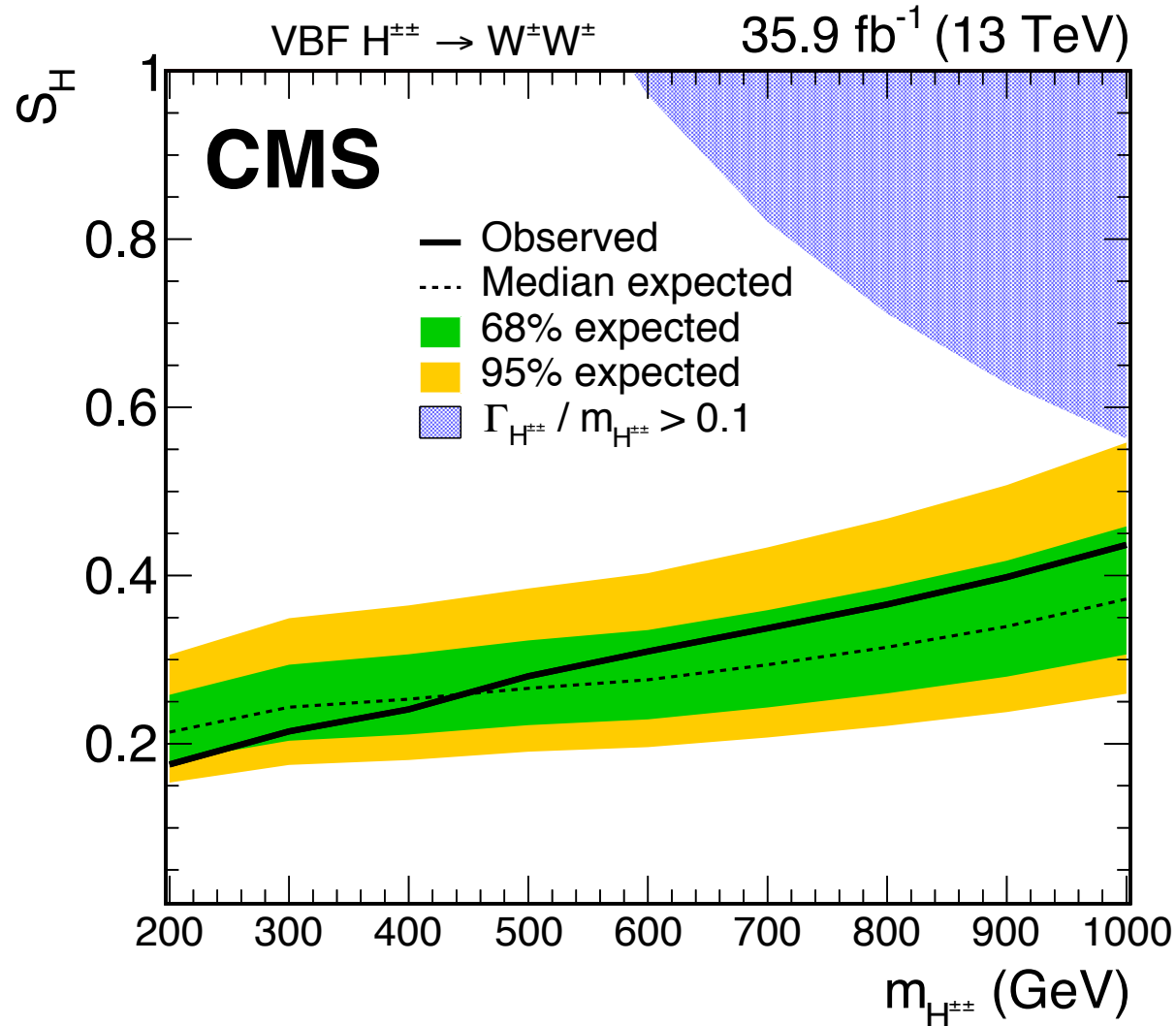
$$\text{VBF} \rightarrow H_5^\pm \rightarrow W^\pm Z$$

VBF +  $qql$ ; VBF +  $3l$  + MET



Cross section  $\propto s_H^2 \equiv$  fraction of  $M_W^2, M_Z^2$  due to exotic scalars

Most stringent constraint: CMS, arXiv:1709.05822



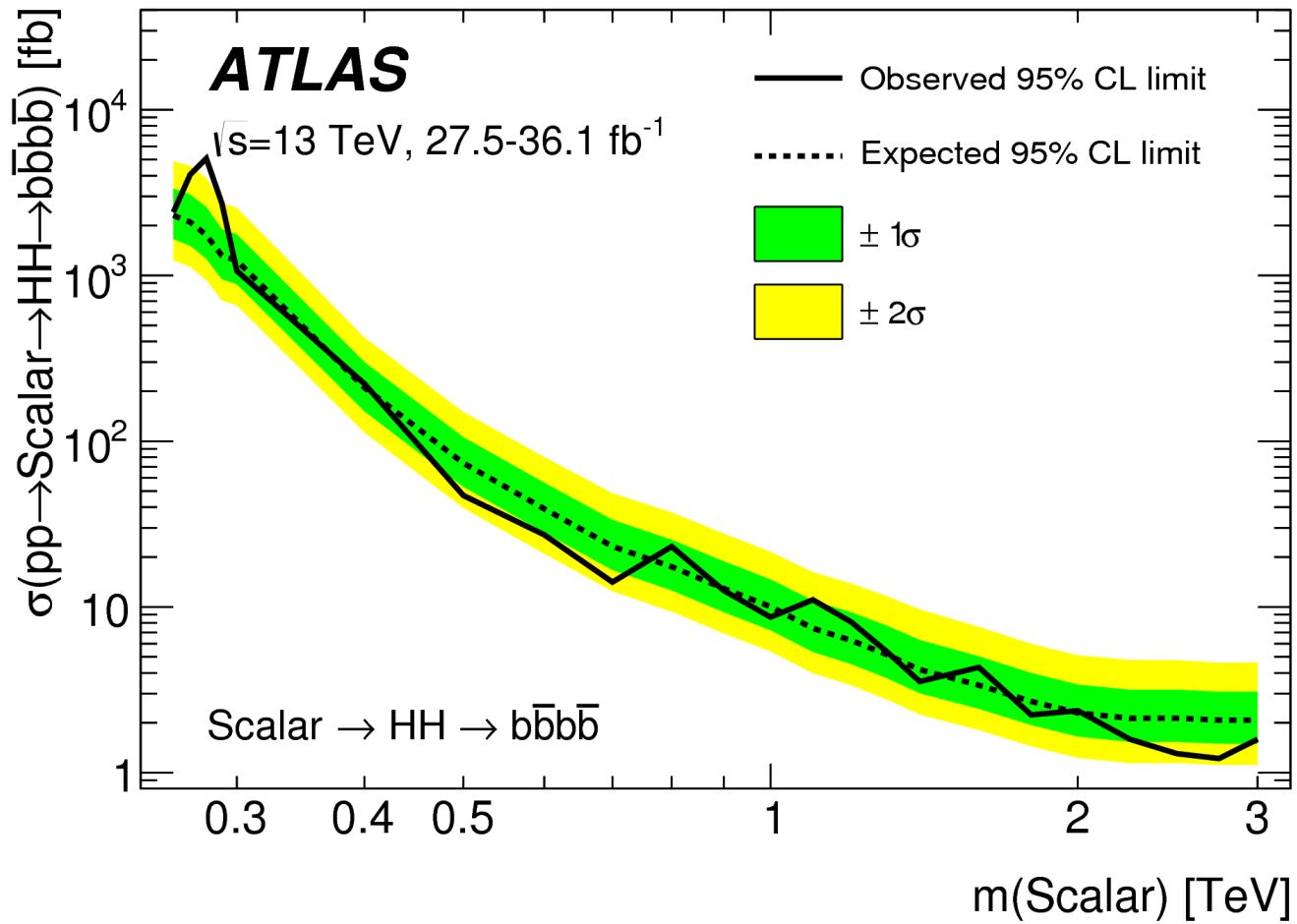
Searches for VBF  $\rightarrow H_5^{\pm} \rightarrow W^{\pm}Z$ : not quite as constraining

## Can GM model provide a benchmark for $H \rightarrow hh$ ?

- Custodial singlet  $H$  couples to fermion pairs  
→ production via gluon fusion
- $Hhh$  coupling can be substantial
- $h$  BRs are generally SM-like in allowed parameter regions

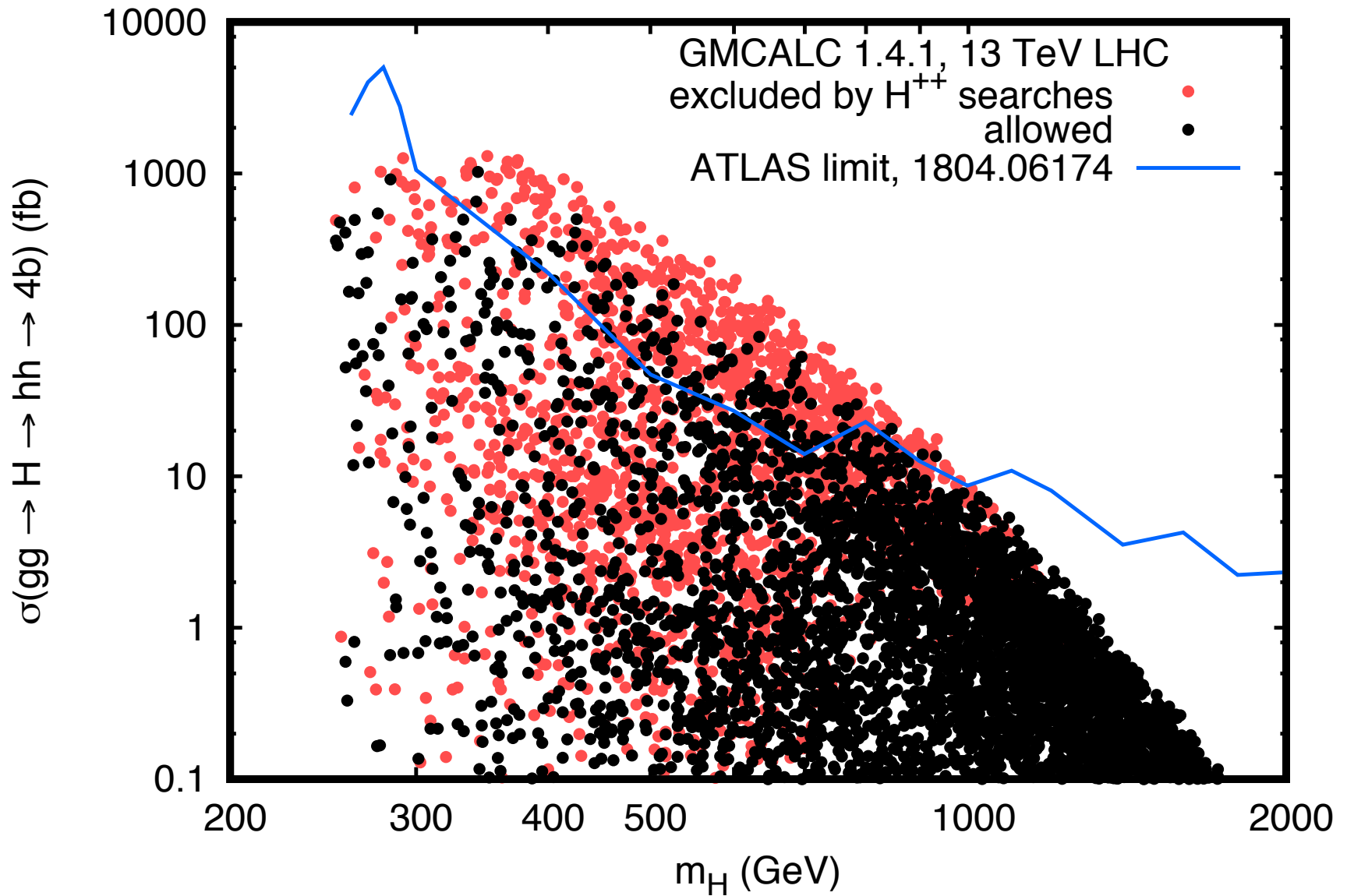
$Hf\bar{f}$  and  $Hhh$  couplings both go to zero in  $s_H \rightarrow 0$  limit, but direct-search constraints are still far away from this limit.

Cross section constrained so far: [ATLAS, arXiv:1804.06174](#)

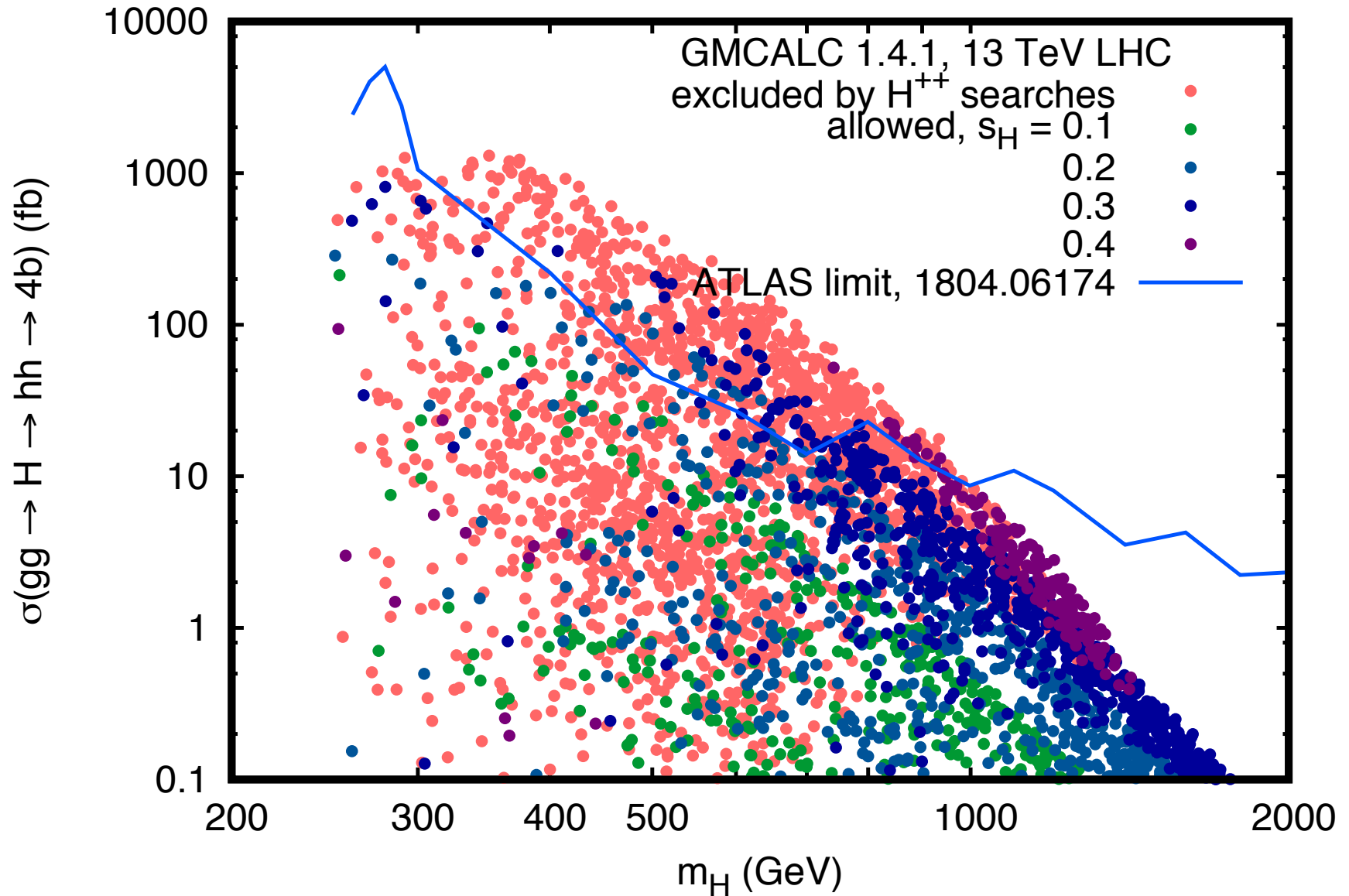




Full parameter scan in GM model: (Preliminary)



Dependence on  $s_H$ : (Preliminary)



## Summary:

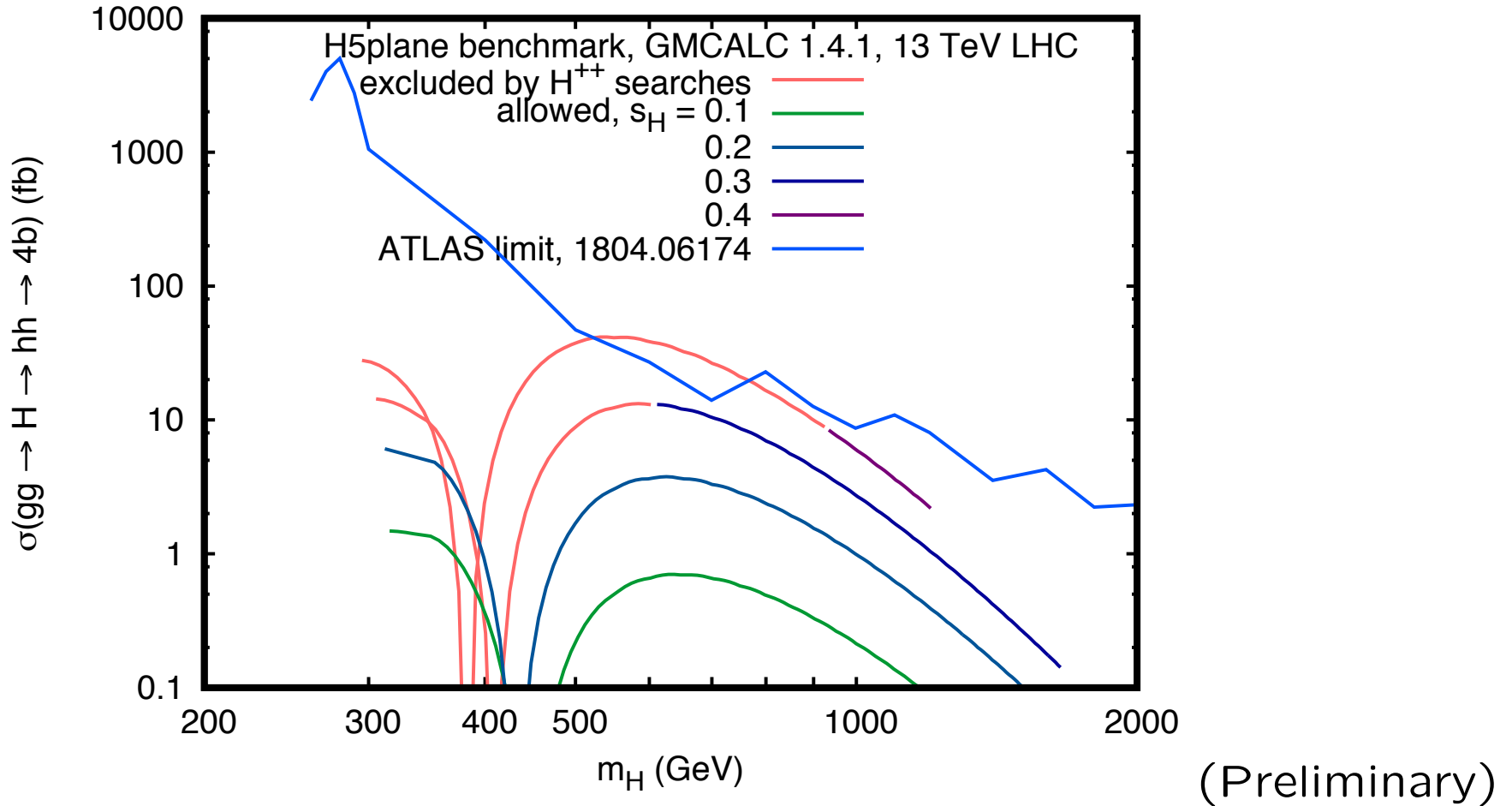
Georgi-Machacek model can provide an interesting benchmark for  $gg \rightarrow H \rightarrow hh$  searches.

Current  $H \rightarrow hh$  sensitivity ( $\lesssim 36 \text{ fb}^{-1}$ , 13 TeV) probes otherwise-unexcluded model points up to  $m_H \sim 1 \text{ TeV}$ !

Feedback wanted: how best to provide this model interpretation to the experiments?

# BACKUP SLIDES

The H5plane benchmark is not so interesting for  $H \rightarrow hh$  searches:



H5plane benchmark designed for  $VBF \rightarrow H_5^{\pm\pm}, H_5^\pm$  searches; two free parameters are  $m_5$  and  $s_H$ , other parameters fixed

Georgi-Machacek model [Georgi & Machacek 1985](#); [Chanowitz & Golden 1985](#)

SM Higgs (bi-)doublet + triplets (1, 0) + (1, 1) in a **bi-triplet**:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix} \quad X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{++*} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

Global  $SU(2)_L \times SU(2)_R \rightarrow$  custodial symmetry  $\langle \chi^0 \rangle = \langle \xi^0 \rangle \equiv v_\chi$

Most general scalar potential invariant under  $SU(2)_L \times SU(2)_R$ :

$$\begin{aligned} V(\Phi, X) = & \frac{\mu_2^2}{2} \text{Tr}(\Phi^\dagger \Phi) + \frac{\mu_3^2}{2} \text{Tr}(X^\dagger X) + \lambda_1 [\text{Tr}(\Phi^\dagger \Phi)]^2 \\ & + \lambda_2 \text{Tr}(\Phi^\dagger \Phi) \text{Tr}(X^\dagger X) + \lambda_3 \text{Tr}(X^\dagger X X^\dagger X) \\ & + \lambda_4 [\text{Tr}(X^\dagger X)]^2 - \lambda_5 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) \text{Tr}(X^\dagger t^a X t^b) \\ & - M_1 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) (UXU^\dagger)_{ab} - M_2 \text{Tr}(X^\dagger t^a X t^b) (UXU^\dagger)_{ab} \end{aligned}$$

9 parameters, 2 fixed by  $G_F$  and  $m_h \rightarrow 7$  free parameters. [Aoki & Kanemura, 0712.4053](#)

[Chiang & Yagyu, 1211.2658](#); [Chiang, Kuo & Yagyu, 1307.7526](#)

[Hartling, Kumar & HEL, 1404.2640](#)

Both approaches have theoretical “issues”:

1) Can't give the septet a vev through spontaneous breaking without generating a physical massless Goldstone boson.

Have to couple it to the SM doublet through a dimension-7  $X\Phi^*\Phi^5$  term [Hisano & Tsumura 2013](#)

Need the UV completion to be nearby!

2) Global  $SU(2)_L \times SU(2)_R$  is broken by gauging hypercharge.

[Gunion, Vega & Wudka 1991](#)

Special relations among params of *full* gauge-invariant scalar potential can only hold at one energy scale: violated by running due to hypercharge. [Garcia-Pepin, Gori, Quiros, Vega, Vega-Morales, Yu 2014](#)

Need the UV completion to be nearby!

This talk: quantify (2) in the Georgi-Machacek model.