

WG3 Extended Scalars* status report

*Formed from the merger of the Charged Higgs and Neutral Extended Scalars subgroups in May 2018

> Heather Logan Carleton University Ottawa, Canada

LHC HXSWG 15th General Assembly Meeting 2018 December 11





No need to search for each one separately – just need to capture the full range of phenomenology so that nothing is missed.

 \rightarrow Identify prototypical signatures and codify as benchmarks

 \rightarrow Prioritize based on how common / universal a signature is across models

 \rightarrow Understand interplay between direct searches and h_{125} coupling measurements to constrain parameter space

Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

Meetings over the past year: (agendas in the backup slides)

- 2018 Mar 13: meeting on signatures for low-mass fermiophobic scalars in Georgi-Machacek model

- Focus on Drell-Yan production of pairs of H_5 states
- $H_5^0 \rightarrow \gamma \gamma$: diphoton resonance fiducial xsec limits \rightarrow recast
- $H_5^{\pm} \rightarrow W^{\pm}\gamma$: study in progress
- $H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$: Run-1 theorist-recast: $m_{H_5^{\pm\pm}} \gtrsim 75$ GeV
- 2018 Sep 20: theory report on $H^{\pm} \rightarrow W^{\pm} \gamma$ to WG3 meeting
 - UFO file available with $H^{\pm}W^{\mp}\gamma$ effective vertex (GM model)

- 2018 Oct 24: open meeting for benchmark proposals & discussion

- 8 talks, many benchmarks; some details in following slides

The simplest extension: SM + real singlet (RxSM)

Two physical Higgs bosons $\phi^{0,r}$ (doublet) and s (singlet)

Mass eigenstates: $m_h < m_H$ $h = \cos \alpha \phi^{0,r} - \sin \alpha s$ $H = \sin \alpha \phi^{0,r} + \cos \alpha s$

All couplings are SM times $\cos \alpha$ for h, SM times $\sin \alpha$ for H. Only possible new decay channel is $H \rightarrow hh$.

 \rightarrow Interpretation for SM-like Higgs coupling measurements as a single overall signal-strength modifier $\mu\equiv\cos\alpha$

 \rightarrow Interpretation for searches for heavy SM-like Higgs boson with overall suppression of all couplings $\mu_H \equiv \sin \alpha$

Combined limits on $|\sin \alpha|$

(A. Ilnicka, TR, T. Stefaniak, Mod.Phys.Lett. A33 (2018) no.10n11, 1830007)



Newest update using Run II results in HiggsSignals: Signal strengths strongest constraint up to 800 GeV: $\sin \alpha \le 0.22$

Heather Logan (Carleton U.) WG3 Extended Scalars

Benchmark for scalar resonance $H \rightarrow hh$ ($h = h_{125}$) in RxSM

Production cross section of $H = \sin^2 \alpha \times \sigma_{SM}(M_H)$

$m_H[{ m GeV}]$	$ \sin lpha _{\sf max}$	$BR_{\min}^{H \to hh}$	$BR_{\max}^{H \rightarrow hh}$	$m_H[GeV]$	$ \sinlpha _{\sf max}$	$BR_{\min}^{H \to hh}$	$BR_{\max}^{H \to hh}$
260	0.22	0.17	0.32	470	0.22	0.23	0.29
270	0.22	0.22	0.37	520	0.22	0.20	0.27
280	0.22	0.23	0.39	590	0.22	0.20	0.26
290	0.22	0.24	0.40	670	0.22	0.20	0.26
310	0.22	0.25	0.40	770	0.22	0.22	0.24
330	0.22	0.25	0.39	880	0.19	0.22	0.25
350	0.22	0.25	0.38	920	0.18	0.22	0.25
370	0.22	0.24	0.36	980	0.17	0.23	0.25
400	0.22	0.22	0.32	1000	0.16	0.23	0.25

Minimal and maximal branching ratios for $H \rightarrow h h$

Tania Robens, WG3 Extended Scalars subgroup meeting, 2018/10/24

Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

Model-specific electroweak radiative corrections calculated!



Real Singlet model

 $H \rightarrow hh$

NLO Corrections shown to be only a few percent

BOJARSKI, CHALONS, LOPEZ-VAL, ROBENS, JHEP1602 (2016) 142

Heather Logan (Carleton U.)

WG3 Extended Scalars

Two Higgs Doublet Model (2HDM)

Physical spectrum: h, H, A, H^{\pm} (assuming CP conservation) Same as in MSSM, but with two key differences:

1) MSSM has "Type II" Yukawa coupling structure; 2HDM can have any of Type I, II, X (Lepton Specific), or Y (Flipped).

2) MSSM quartic scalar couplings are fixed by g, g'; not true in 2HDM: can have much larger mass splittings among H, A, H^{\pm}

Less constrained spectra allow for Higgs-to-Higgs decays: $A \rightarrow ZH / H \rightarrow ZA$ (also $A \rightarrow Zh_{125}$) \rightarrow ATLAS + CMS $H^+ \rightarrow W^+S$ ($S = H, A, h_{125}$) \rightarrow CMS, $S = A(\rightarrow \mu\mu)$ $H \rightarrow AA \rightarrow$ ATLAS, $AA \rightarrow 4\gamma$

 $H \to H^+H^- \to \tau \nu \to \text{not being done}$ $H/A \to W^{\pm}H^{\mp} \to \text{not being done}$

Heather Logan (Carleton U.) WG3 Extended Scalars

"Standard" high-mass charged Higgs search $pp \rightarrow tH^- \rightarrow tW^-b\overline{b}$:

In 2HDM can decay via $H^- \rightarrow W^-A(\rightarrow b\overline{b})$; large H^- and Awidths give large interference between signal and background. A. Arhrib et al., 1712.05018

Study in progress to see how big an issue this is in MSSM:

D. Azevedo, R. Santos, S. Moretti, P. Sharma, R. Benbrik, A. Arhrib & R. Patrick H^+ width can be large enough to lead to significant interference



Heather Logan (Carleton U.)

WG3 Extended Scalars

²⁰¹⁸ Dec 11

Interference effects in H^{\pm} production at the LHC

D. Azevedo, R. Santos, S. Moretti, P. Sharma, Rachid Benbrik, A. Abdesslam, R. Patrick

Benchmarks in

MSSM

Kinematical Distributions

Conclusions

Benchmark cross-sections

 \rightarrow Benchmarks points are chosen where there is a large charged Higgs width and smallest m_A^0 :

Parameters	hMSSM	m_{h}^{mod+}	${\sf m}_{\sf h}^{{125}}(ilde{ au})$
μ (GeV)	200	200	1000
aneta	1.01	3.42	3.19
m_{H^+} (GeV)	633.91	303.08	628.08
Γ_{H^+} (GeV)	27.777	0.925	2.677

Production cross-sections:

Benchmark	Signal (pb)	Background (pb)
hMSSM	$(3.2402\pm0.0014) imes10^{-2}$	13.092 ± 0.004
$m_h^{ m mod+}$	$(8.8502 \pm 0.0033) imes 10^{-2}$	13.103 ± 0.004
$m_h^{125}(ilde{ au})$	$(1.6802\pm0.0058) imes10^{-2}$	13.177 ± 0.004

Benchmark	Signal+Background (pb)	Interference (pb)
hMSSM	13.143 ± 0.004	0.019 ± 0.008
$m_h^{ m mod+}$	13.200 ± 0.004	0.009 ± 0.008
$m_h^{125}(ilde{ au})$	13.197 ± 0.004	0.003 ± 0.008

Where

$$(S+B)^2 = S^2 + B^2 + \text{Interference}$$
(1)

 \rightarrow Still large errors but interferences seem to be present.

Work in progress.

Heather Logan (Carleton U.)

WG3 Extended Scalars

2018 Dec 11

10

2HDM with explicit CP violation (C2HDM)

Physical spectrum: H_1 , H_2 , H_3 , H^{\pm}

3 neutral scalars H_1 , H_2 , H_3 are CP admixtures in general

- Motivated by need for new sources of CP violation to explain baryon asymmetry of the universe

- Constrained by null searches for electric dipole moments

New processes not present in Real 2HDM:

- $H \rightarrow SS \rightarrow 4W \ (S \neq h_{125})$: $\rightarrow \text{ATLAS}$

- $H_3 \rightarrow H_1 H_2$: one of these must be h_{125} ; motivates $H \rightarrow h_{125}S$ selection ($m_S \neq 125 \text{ GeV}$) \rightarrow not being done

Both of these can also happen in CP-conserving 2HDM + real singlet ("N2HDM"), which has 3 CP-even neutral Higgs bosons.

Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

What do we know about possible CP-violating couplings of h_{125} ?

CP properties have been tested so far in hZZ, hWW couplings - $h \rightarrow 4\ell$ distributions, production distributions in VBF

CP-even SM $hV_{\mu}V^{\mu}$ is tree level / dim-4 operator CP-odd $hV_{\mu\nu}\tilde{V}^{\mu\nu}$ is one-loop / dim-6 operator \rightarrow CP-odd coupling in hVV is generically small

Next place to look is the Yukawa couplings. For Type-II Yukawas:

t:
$$Y_{C2HDM}^{TypeII} = \cos \alpha_2 Y_{2HDM}^{TypeII} - i\gamma_5 \sin \alpha_2 \cot \beta$$

b, τ : $Y_{C2HDM}^{TypeII} = \cos \alpha_2 Y_{2HDM}^{TypeII} - i\gamma_5 \sin \alpha_2 \tan \beta$

 α_2 is mixing angle between pseudoscalar and scalars; $\kappa_V = \cos \alpha_2 \sin(\beta - \alpha)$ so $|\cos \alpha_2|$ must be near 1 already.

But, $\tan \beta$ can be large: look for CP-violating effects in the Yukawas with $\tan \beta$ enhancement!

Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

Rate measurements constrain CP-even and CP-odd parts of Yukawa couplings to lie in a ring: $Y = a + ib\gamma_5$, rate $\propto |a|^2 + |b|^2$

Electric dipole moment measurements (of n, atoms, molecules) constrain the amount of CP violation.

Depends on Yukawa structure!

Large CPV Yukawas excluded in Type I and

in Type II when $h_{125} =$ lightest neutral scalar.

Most interesting scenarios:

Fontes et al, 1711.09419



Heather Logan (Carleton U.)

WG3 Extended Scalars



Probe using τ decay distributions! Benchmark points available. Fontes et al, 1711.09419

Heather Logan (Carleton U.)

WG3 Extended Scalars

Higgs Triplet Model (HTM)

SM Higgs doublet plus 1 complex triplet $X = (\chi^{++}, \chi^{+}, \chi^{0})$

Motivation: $y_{\nu}^{ij}L_iXL_j$ coupling gives neutrino masses $m_{\nu} \sim y_{\nu} \langle \chi^0 \rangle$

 $\langle \chi^0 \rangle$ is very strongly constrained by the ρ parameter:

$$\rho \equiv \frac{\text{weak neutral current}}{\text{weak charged current}} = \frac{(g^2 + g'^2)/M_Z^2}{g^2/M_W^2} = \frac{v_\phi^2 + a\langle X^0 \rangle^2}{v_\phi^2 + b\langle X^0 \rangle^2}$$

$$a = 4 \left[T(T+1) - Y^2 \right] c$$

$$b = 8Y^2$$

$$Q = T^3 + Y; \text{ SM doublet: } Y = 1/2$$

Expt: $\rho = 1.00039 \pm 0.00019$ (2018 PDG)

 \Rightarrow $\langle \chi^0 \rangle \lesssim$ GeV; negligible mixing of χ^0 with SM-like Higgs

Heather Logan (Carleton U.) WG3 Extended Scalars

 $\chi^{\pm\pm}$ decays to $\ell^\pm\ell^\pm$ or $W^\pm W^\pm$ depending on size of triplet vev



Rely on Drell-Yan production of $\chi^{++}\chi^{--}$ or $\chi^{\pm\pm}\chi^{\mp}$.

Like-sign dilepton resonace search is very sensitive – exclude $\chi^{\pm\pm}$ up to ~ 800 GeV depending on assumptions about $e/\mu/\tau$ fractions \rightarrow ATLAS + CMS

Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

 $\chi^{\pm\pm} \to W^{\pm}W^{\pm}$ search done for first time in Run 2 (Ws on shell)



Theorist recast of ATLAS Run-1 like-sign dimuon data sets lower bound $m_{\chi^{++}}\gtrsim$ 84 GeV Kanemura, Kikuchi, Yagyu & Yokoya, 1412.7603

Gap at intermediate masses< 200 GeV: need offshell Ws!</th>Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

17

Models with triplets (or larger) contributing to EWSB:

Have to model-build to avoid $\rho \neq 1$. Only two known approaches:

1) Use the septet (T, Y) = (3, 2): $\rho = 1$ by accident! Doublet $(\frac{1}{2}, \frac{1}{2})$ + septet (3, 2): Scalar septet model Hisano & Tsumura, 1301.6455; Kanemura, Kikuchi & Yagyu, 1301.7303

2) Use global $SU(2)_L \times SU(2)_R$ imposed on the scalar potential Global $SU(2)_L \times SU(2)_R \rightarrow$ custodial SU(2) ensures tree-level $\rho = 1$ Doublet + triplets (1,0) + (1,1): Georgi-Machacek model

Georgi & Machacek 1985; Chanowitz & Golden 1985Doublet + quartets $(\frac{3}{2}, \frac{1}{2}) + (\frac{3}{2}, \frac{3}{2})$:Generalized Georgi-
Doublet + quintets (2, 0) + (2, 1) + (2, 2):Doublet + quintets $(\frac{5}{2}, \frac{1}{2}) + (\frac{5}{2}, \frac{3}{2}) + (\frac{5}{2}, \frac{5}{2})$:Machacek models

Galison 1984; Robinett 1985; HEL 1999; Chang et al 2012; HEL & Rentala 2015 Larger than sextets \rightarrow too many large multiplets, violates perturbativity

<u>All</u> contain doubly-charged Higgs with $H^{\pm\pm}W^{\mp}W^{\mp}$ coup. $\propto \langle X^0 \rangle$!

Heather Logan (Carleton U.) WG3 Extended Scalars

Georgi-Machacek model Georgi & Machacek 1985; Chanowitz & Golden 1985

SM Higgs (bi-)doublet + triplets (1,0) + (1,1) in a bi-triplet:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix} \qquad X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{++*} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

Global SU(2)_L×SU(2)_R \rightarrow custodial symmetry $\langle \chi^0 \rangle = \langle \xi^0 \rangle \equiv v_{\chi}$

Physical spectrum: Bi-doublet: $2 \otimes 2 \rightarrow 1 \oplus 3$

 $\mathsf{Bi-triplet:} \ \mathbf{3}\otimes\mathbf{3}\to\mathbf{1}\oplus\mathbf{3}\oplus\mathbf{5}$

- Two custodial singlets mix $\rightarrow h$, $H \ m_h$, m_H , angle α Usually identify h = h(125)
- Two custodial triplets mix $\rightarrow (H_3^+, H_3^0, H_3^-) m_3 + \text{Goldstones}$ Phenomenology very similar to H^{\pm}, A^0 in 2HDM Type I, $\tan \beta \rightarrow \cot \theta_H$
- Custodial fiveplet $(H_5^{++}, H_5^+, H_5^0, H_5^-, H_5^{--}) m_5$ Fermiophobic; H_5VV couplings $\propto s_H \equiv \sqrt{8}v_{\chi}/v_{\rm SM}$ $s_H^2 \equiv$ exotic fraction of M_W^2 , M_Z^2

Heather Logan (Carleton U.) WG3 Extended Scalars

Explicit LHC searches up to now:

$$VBF \rightarrow H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm} \rightarrow CMS \quad VBF + \text{like-sign dileptons} + \text{MET}$$
$$VBF \rightarrow H_5^{\pm} \rightarrow W^{\pm}Z \rightarrow \text{ATLAS} + CMS \quad VBF \ qq\ell\ell; \ VBF \ 3\ell + \text{MET}$$



Cross section $\propto s_{H}^{2} \equiv$ fraction of M_{W}^{2}, M_{Z}^{2} due to exotic scalars

Heather Logan (Carleton U.)

WG3 Extended Scalars

Most stringent constraint: VBF $\rightarrow H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ CMS, arXiv:1709.05822



Also searches for VBF $\rightarrow H_5^{\pm} \rightarrow W^{\pm}Z \rightarrow \text{ATLAS} + \text{CMS}$ Heather Logan (Carleton U.) WG3 Extended Scalars 2018 Dec 11



Recast ATLAS Run1 $\gamma\gamma$ resonance, GMCALC 1.5.0 beta

Extending Drell-Yan $H^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ search to masses below 200 GeV (w/ offshell Ws) could exclude entire low- m_5 region!

Heather Logan (Carleton U.)

WG3 Extended Scalars

For m_5 below threshold for $H_5 \rightarrow VV$ (V = W, Z) decays, BRs for loop-induced decays $H_5^0 \rightarrow \gamma\gamma$, $H_5^{\pm} \rightarrow W^{\pm}\gamma$ become important (remember H_5 is fermiophobic)

Logan & Wu, 1809.09127 Sensitivity study for $H_5^{\pm} \rightarrow W^{\pm}\gamma$ (production by Drell-Yan):



Expected 95%CL exclusion with 300 fb⁻¹ at 14 TeV LHC Recast ATLAS Run1 $\gamma\gamma$ resonance search current exclusion

 $H_5^{\pm} \rightarrow W^{\pm}\gamma$ simulation tool now public: UFO model for MG5

If low- m_5 region in GM model is excluded by $H^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$, anew theory benchmark to motivate $H^+ \rightarrow W^+\gamma$ will be needed.Heather Logan (Carleton U.)WG3 Extended Scalars2018 Dec 11

Several recent "low-mass" results rely on Drell-Yan production of pairs of new scalars.

Request to provide Drell-Yan xsec tables (in progress, NLO QCD)

- GM model $pp \to H_5^{++}H_5^{--}$, $H_5^{\pm\pm}H_5^{\mp}$, $H_5^{+}H_5^{-}$, $H_5^{\pm}H_5^{0}$

- HTM
$$pp \rightarrow \chi^{++}\chi^{--}$$
, $\chi^{\pm\pm}\chi^{\mp}$ etc.

- 2HDM

Simple relations between cross sections in different models due to gauge quantum numbers of scalars.

GM model benchmark for $H \rightarrow hh$: full parameter scan (Prelim)





WG3 Extended Scalars

²⁰¹⁸ Dec 11

Summary: "wish list" of missing search channels

 $H_3 \rightarrow H_1 H_2$, where all three Higgs bosons have different masses. h_{125} could be any of these three. (our #1 priority for a new search)

 $H^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$: extend search to masses below 200 GeV (off-shell Ws). Production via Drell-Yan in pairs or with H^{\mp} .

$$H \to H^+ H^- \to \tau \nu \tau \nu \ (H \neq h_{125})$$

$$H \to W^+ H^- \ (H \neq h_{125})$$

 $H^+ \rightarrow W^+ \gamma$: search for fermiophobic charged Higgs including at low mass (below 200 GeV); production via Drell-Yan.

 $H_{125} \rightarrow \tau \tau$ CP measurement from τ polarization kinematic distributions [this belongs to SM Higgs Characterization]

Heather Logan (Carleton U.) WG3 Extended Scalars

the end

Heather Logan (Carleton U.)

WG3 Extended Scalars

BACKUP SLIDES

Heather Logan (Carleton U.)

WG3 Extended Scalars

2018 Dec 11

28

WG3: Extended Scalars meeting					
\blacksquare Tuesday 13 Mar 2018, 14:00 \rightarrow 16:00 Europe/Zurich					
Vidyo or	♥ Vidvo only (CERN)				
Videoconferer Roo	ms WG3_Extended_Scalars_meeting	Join 😽			
There are	minutes attached to this event. Show them.				
14:00 → 14:05	Introduction Speakers: Heather Logan (Carleton University), Dr Raffaele Angelo Gerosa (Univ. of California San Diego (US)), Rui Santos (IST), Shufang Su (University of Arizona), Xiangyang Ju (University of Wisconsin Madison (US)), Xiaohu Sun (University of Alberta (CA))	③ 5m			
14:05 → 14:25	Drell-Yan H5^0> gamma gamma Speaker: Roberto Vega-Morales	() 20m			
14:25 → 14:45	H5+> W+ gamma theory Speaker: Yongcheng Wu IMSWGAinGM_ycwu	() 20m			
14:45 → 15:05	H5+> W+ gamma experiment Speakers: Brigitte Vachon, Kays Haddad, Kays Haddad (McGill University, (CA)) Charged Higgs Grou	() 20m			
15:05 → 15:25	Drell-Yan H5 ⁺⁺ > W+ W+> like-sign leptons Speaker: Heather Logan (Carleton University) M5pp-to-leptons.pdf	() 20m			
15:25 → 15:35	Discussion	🕲 10m			



WG3: Benchmark Discussion Wednesday 24 Oct 2018, 17:00 → 20:00 Europe/Zurich			
Videoconfere Roc	ms WG3_Benchmark_Discussion	Join 🔶	
17:00 → 17:05	Introduction Speaker: Rui Santos (ISEL and CFTC-UL)	© 5m	
17:05 → 17:15	Benchmark Points for Type-I 2HDMs with a light h Speaker: William Klemm	③ 10m	
17:15 → 17:25	Benchmark Point with low-mass CP-odd Higgs A with strong couplings to leptons and top quarks Speaker: Dominik Stoeckinger	© 10m	
17:25 → 17:35	Charged Higgs boson benchmarks from top quark polarization Speaker: Adil Jueid	③ 10m	
17:35 → 17:45	IDM benchmarks for the LHC at 13 and 27 TeV Speaker: Tania Robens	③10m	
17:45 → 17:55	Updated constraints for the Real Higgs Singlet Extension of the Standard Model Speaker: Tania Robens Singlet_robens.pdf	③ 10m	
17:55 → 18:05	Interference effects in H± production at the LHC Speaker: Duarte Azevedo	© 10m	
18:05 → 18:15	A benchmark for LHC searches for H5±±, H5±, and H05 in the Georgi-Machacek model including masses below 200 GeV Speaker: Ben Keeshan BenKeeshanHXSW	③ 10m	
18:15 → 18:25	Benchmark scenarios in the C2HDM Speaker: Jonas Wittbrodt	③ 10m	
18:25 → 18:45	Discussion	③ 20m	

 $H \to hh~{\rm cross~section~constrained~so~far:}~{\rm ATLAS,~arXiv:1804.06174}$



Georgi-Machacek model Georgi & Machacek 1985; Chanowitz & Golden 1985

SM Higgs (bi-)doublet + triplets (1,0) + (1,1) in a bi-triplet:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix} \qquad X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{++*} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

Global SU(2)_L×SU(2)_R \rightarrow custodial symmetry $\langle \chi^0 \rangle = \langle \xi^0 \rangle \equiv v_{\chi}$

Most general scalar potential invariant under $SU(2)_L \times SU(2)_R$:

$$V(\Phi, X) = \frac{\mu_2^2}{2} \operatorname{Tr}(\Phi^{\dagger} \Phi) + \frac{\mu_3^2}{2} \operatorname{Tr}(X^{\dagger} X) + \lambda_1 [\operatorname{Tr}(\Phi^{\dagger} \Phi)]^2 + \lambda_2 \operatorname{Tr}(\Phi^{\dagger} \Phi) \operatorname{Tr}(X^{\dagger} X) + \lambda_3 \operatorname{Tr}(X^{\dagger} X X^{\dagger} X) + \lambda_4 [\operatorname{Tr}(X^{\dagger} X)]^2 - \lambda_5 \operatorname{Tr}(\Phi^{\dagger} \tau^a \Phi \tau^b) \operatorname{Tr}(X^{\dagger} t^a X t^b) - M_1 \operatorname{Tr}(\Phi^{\dagger} \tau^a \Phi \tau^b) (U X U^{\dagger})_{ab} - M_2 \operatorname{Tr}(X^{\dagger} t^a X t^b) (U X U^{\dagger})_{ab}$$

9 parameters, 2 fixed by G_F and $m_h \rightarrow 7$ free parameters. Aoki & Kanemura, 0712.4053 Chiang & Yagyu, 1211.2658; Chiang, Kuo & Yagyu, 1307.7526 Hartling, Kumar & HEL, 1404.2640

WG3 Extended Scalars 2018 Dec 11

Heather Logan (Carleton U.)

 $H^{++} \rightarrow W^+W^+$: Below the WW threshold, same-flavour leptons can come from either of the Ws, leading to an interference term. Need full $H^{++} \rightarrow 4f$ branching ratios simulation.



Kanemura, Kikuchi, Yagyu & Yokoya, 1407.6547

Heather Logan (Carleton U.)

WG3 Extended Scalars

HEL & Rentala, "All the generalized Georgi-Machacek models," 1502.01275



Set limit $m_{H^{++}} \gtrsim$ 76 GeV in GM model.

Heather Logan (Carleton U.)

WG3 Extended Scalars