



Searching for the $W\gamma$ decay of a charged Higgs boson

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Prospects for Charged Higgs Discovery at Colliders
(CHARGED 2018)
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Based on [HEL & Yongcheng Wu, arXiv:1809.09127](#)



Outline

Introduction

$H^\pm \rightarrow W^\pm \gamma$ calculation and effective vertex

Concrete implementation: Georgi-Machacek model

LHC collider study & results

Conclusions

Introduction

Why charged Higgs:



Charged Higgs searches up to now:

Fermiophilic:

- Production via $t\bar{b}H^+$ coupling, decays to fermions
- Production in decays of heavier Higgs (2HDM $H/A \rightarrow H^+W^-$)
(proposed)

Fermiophobic:

- VBF production of H^\pm , decays to $W^\pm Z$

This talk:

- Loop-induced decay $H^\pm \rightarrow W^\pm \gamma$

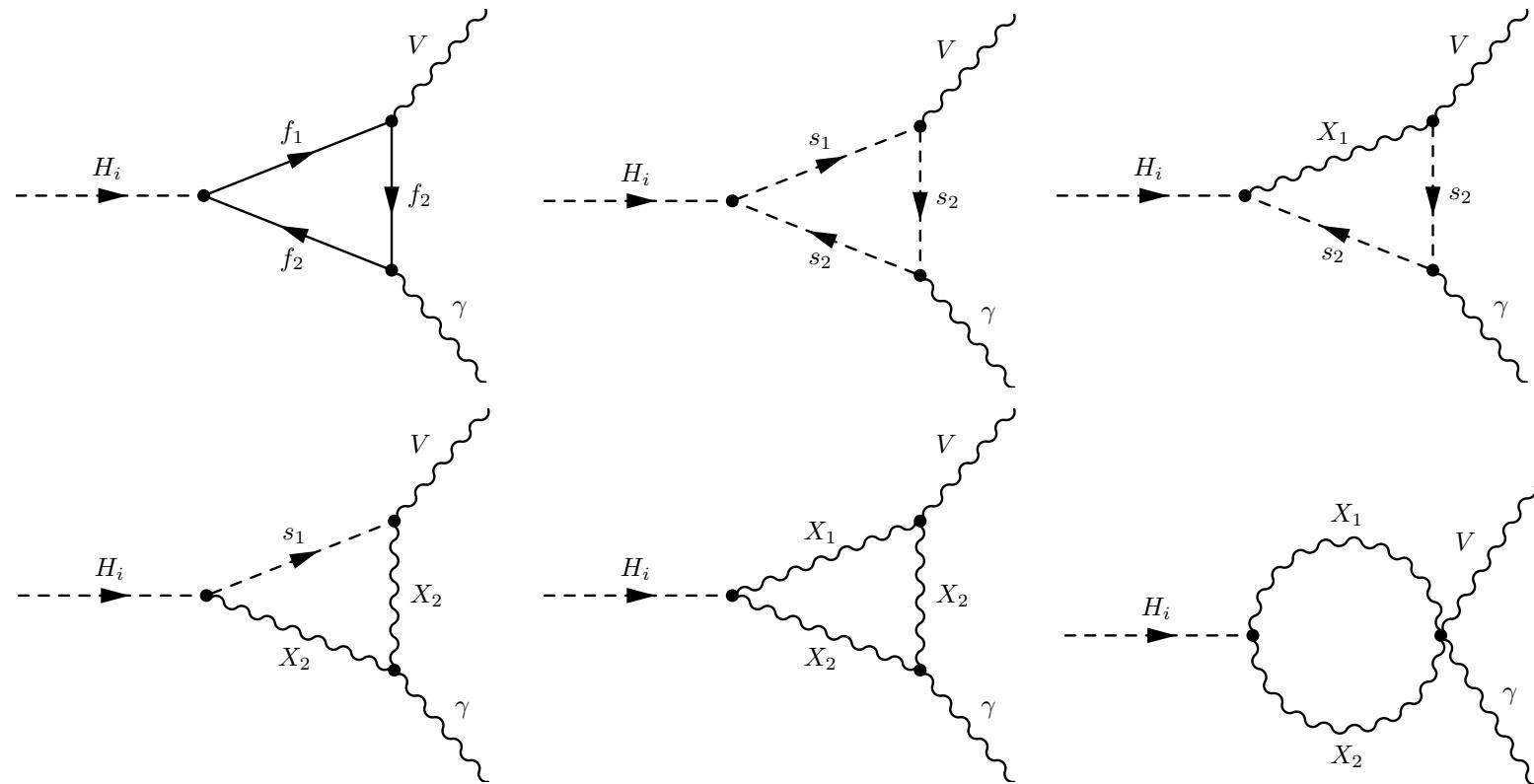
Loop calculations for $H^+ \rightarrow W^+ \gamma$

2HDM/MSSM: Arhrib, Benbrik & Chabab, hep-ph/0607182

Stealth Doublet model: Enberg, Rathsman & Wouda, 1311.4367

Aligned 2HDM: Ilisie & Pich, 1405.6639

Georgi-Machacek model: Degrande, Hartling & HEL, 1708.08753



Most interesting when H^\pm is **fermiophobic** and **light** ($< M_W + M_Z$)

One-loop effective vertex

2 form factors: S (scalar) and \tilde{S} (pseudoscalar)

$$\mathcal{M} = \Gamma^{\mu\nu} \varepsilon_\nu^{W*}(k) \varepsilon_\mu^{\gamma*}(q), \quad \Gamma^{\mu\nu} = (g^{\mu\nu} k \cdot q - k^\mu q^\nu) S + i\epsilon^{\mu\nu\alpha\beta} k_\alpha q_\beta \tilde{S}$$

In a CP-conserving theory:

S : from fermion, scalar, gauge, & mixed gauge/scalar loops

\tilde{S} : only from fermion loop **Fermiophobic $H^\pm \Rightarrow \tilde{S} = 0$**

Decay width:

$$\Gamma(H^+ \rightarrow W^+ \gamma) = \frac{m_{H^+}^3}{32\pi} \left[1 - \frac{m_W^2}{m_{H^+}^2} \right]^3 [|S|^2 + |\tilde{S}|^2],$$

Experimentally, we'll look for $W \rightarrow \ell\nu$.

⇒ Information about W polarization in angular distribution of ℓ relative to γ

One-loop effective vertex & kinematic shapes

2 form factors: S (scalar) and \tilde{S} (pseudoscalar)

$$\mathcal{M} = \Gamma^{\mu\nu} \varepsilon_\nu^{W*}(k) \varepsilon_\mu^{\gamma*}(q), \quad \Gamma^{\mu\nu} = (g^{\mu\nu} k \cdot q - k^\mu q^\nu) S + i\epsilon^{\mu\nu\alpha\beta} k_\alpha q_\beta \tilde{S}$$

Differential distribution:

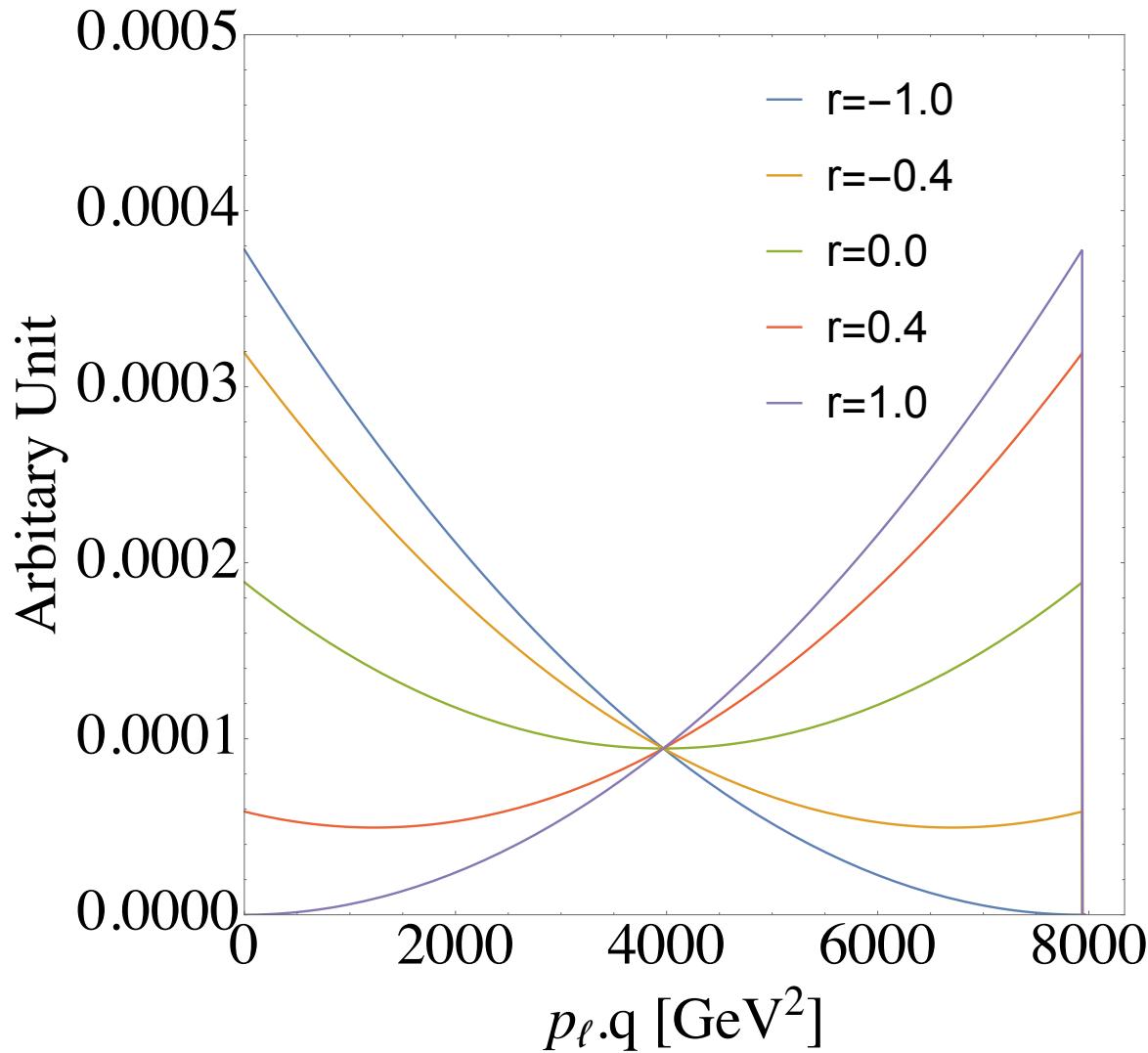
$$\begin{aligned} |\mathcal{M}|^2 &\propto \Gamma^{\mu\nu} \Gamma^{\rho\sigma*} \varepsilon_\mu^{\gamma*} \varepsilon_\rho^\gamma \mathbf{Tr}(\not{p}_\nu \gamma_\sigma P_L \not{p}_\ell \gamma_\nu) \\ &= \frac{m_W^2}{2} \left\{ 8(p_\ell \cdot q)^2 [|S|^2 + |\tilde{S}|^2] \right. \\ &\quad - 4(p_\ell \cdot q)(m_{H^+}^2 - m_W^2) [|S|^2 + |\tilde{S}|^2 - 2\mathbf{Re}(S\tilde{S}^*)] \\ &\quad \left. + (m_{H^+}^2 - m_W^2)^2 [|S|^2 + |\tilde{S}|^2 - 2\mathbf{Re}(S\tilde{S}^*)] \right\} \end{aligned}$$

Define $r \equiv \frac{\tilde{S}}{S}$, $K \equiv \frac{p_\ell \cdot q}{(m_{H^+}^2 - m_W^2)/2} \in [0, 1]$

$$|\mathcal{M}|^2 \propto 2K^2 [1 + |r|^2] + (-2K + 1) [1 + |r|^2 - 2\mathbf{Re}(r)]$$

Parabola in $K \propto \not{p}_\ell \cdot \not{q}$, shape depends on r

One-loop effective vertex & kinematic shapes



$m_{H^+} = 150 \text{ GeV}$

(changing m_{H^+}
just rescales x axis)

Fermiophobic H^+
 $\Rightarrow r = 0$

Concrete implementation: Georgi-Machacek model

Georgi & Machacek 1985; Chanowitz & Golden 1985

SM Higgs (bi-)doublet + triplets $(1, 0) + (1, 1)$ in a bi-triplet:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix} \quad X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{++*} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

Global $SU(2)_L \times SU(2)_R \rightarrow$ custodial symmetry $\langle \chi^0 \rangle = \langle \xi^0 \rangle \equiv v_\chi$
 (ensures $\rho = 1$)

Most general scalar potential invariant under $SU(2)_L \times SU(2)_R$:

$$\begin{aligned} V(\Phi, X) = & \frac{\mu_2^2}{2} \text{Tr}(\Phi^\dagger \Phi) + \frac{\mu_3^2}{2} \text{Tr}(X^\dagger X) + \lambda_1 [\text{Tr}(\Phi^\dagger \Phi)]^2 \\ & + \lambda_2 \text{Tr}(\Phi^\dagger \Phi) \text{Tr}(X^\dagger X) + \lambda_3 \text{Tr}(X^\dagger X X^\dagger X) \\ & + \lambda_4 [\text{Tr}(X^\dagger X)]^2 - \lambda_5 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) \text{Tr}(X^\dagger t^a X t^b) \\ & - M_1 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) (UXU^\dagger)_{ab} - M_2 \text{Tr}(X^\dagger t^a X t^b) (UXU^\dagger)_{ab} \end{aligned}$$

9 parameters, 2 fixed by G_F and $m_h \rightarrow$ 7 free parameters. Aoki & Kanemura, 0712.4053

Chiang & Yagyu, 1211.2658; Chiang, Kuo & Yagyu, 1307.7526

Hartling, Kumar & HEL, 1404.2640

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Physical spectrum:

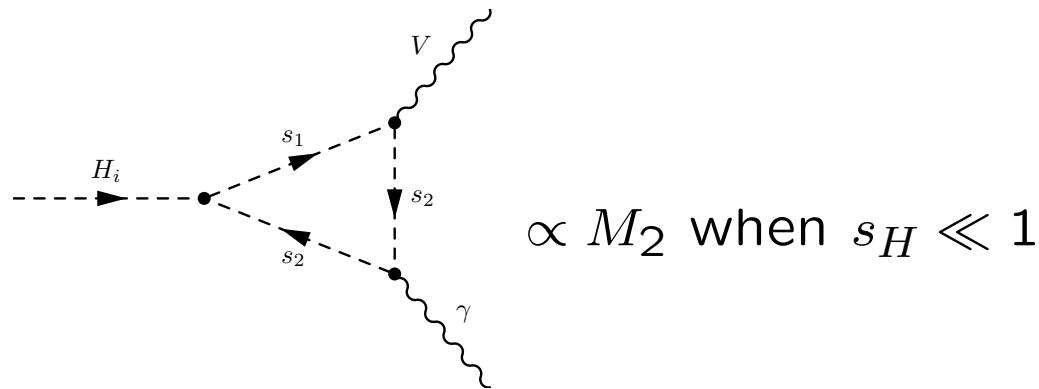
Bi-doublet: $2 \otimes 2 \rightarrow 1 \oplus 3$

Bi-triplet: $3 \otimes 3 \rightarrow 1 \oplus 3 \oplus 5$

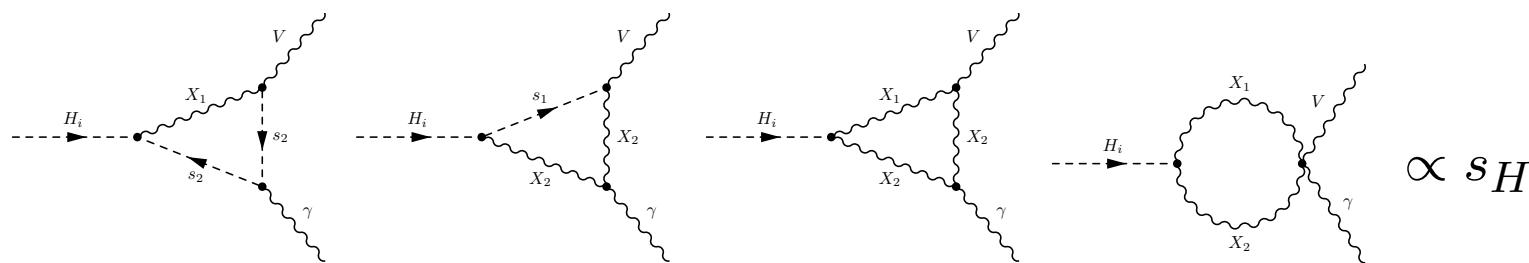
- Two custodial singlets mix $\rightarrow h^0, H^0$ m_h, m_H , angle α
Usually identify $h^0 = h(125)$
- Two custodial triplets mix $\rightarrow (H_3^+, H_3^0, H_3^-)$ m_3 + Goldstones
Phenomenology very similar to H^\pm, A^0 in 2HDM Type I, $\tan \beta \rightarrow \cot \theta_H$
- Custodial fiveplet $(H_5^{++}, H_5^+, H_5^0, H_5^-, H_5^{--})$ m_5
Fermiophobic! $H_5 VV$ couplings $\propto s_H \equiv \sqrt{8}v_\chi/v_{\text{SM}}$
 $s_H^2 \equiv$ exotic fraction of M_W^2, M_Z^2

$H_5^+ \rightarrow W^+ \gamma$ in Georgi-Machacek model

Degrande, Hartling & HEL, 1708.08753



$$\propto M_2 \text{ when } s_H \ll 1$$



Competing decay: $H_5^+ \rightarrow W^+ Z$ coupling $\propto s_H$.

So the most interesting model parameters are m_5 , s_H , and M_2 .

Choice of benchmark

9 Lagrangian parameters: $\mu_2^2, \mu_3^2, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5, M_1, M_2$

Can re-express as: $G_F, m_h, m_H, m_3, m_5, s_H, \sin \alpha, M_1, M_2$

- Fix G_F and $m_h = 125$ GeV
- Keep m_5, s_H , and M_2 as variable parameters
- Fix the remaining 4 parameters as follows:

$$\begin{aligned}m_3^2 &= m_5^2 + \delta m^2 & \delta m^2 &= (300 \text{ GeV})^2 \\m_H^2 &= m_5^2 + \frac{3}{2}\delta m^2 + \kappa_H v^2 s_H^2 & \kappa_H &= -m_5/(100 \text{ GeV}) \\M_1 &= s_H \left[\frac{\sqrt{2}}{v} \left(m_5^2 + \frac{3}{2}\delta m^2 \right) + 3M_2 s_H + \kappa_{\lambda_3} v s_H^2 \right] & \kappa_{\lambda_3} &= -\kappa_H^2/10 \\\sin \alpha &= \kappa_\alpha s_H & \kappa_\alpha &= -0.15 - m_5/(1000 \text{ GeV})\end{aligned}$$

Sufficient for this study: fully populates low m_5 and $s_H < 0.3$.
(Consistency with indirect & h_{125} constraints not checked yet.)

Computational tools

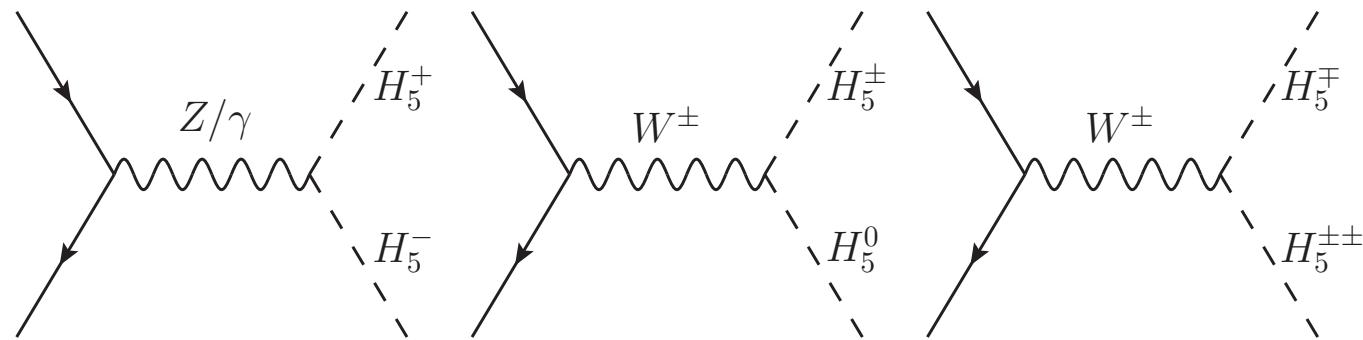
Loop-induced decays for GM model calculated in [Degrade, Hartling & HEL, 1708.08753](#) and implemented in GMCALC 1.3.0 (uses LoopTools)

FeynRules implementation + UFO model file updated to include loop-induced scalar-gauge-gauge vertices in EFT approach (gg , $\gamma\gamma$, $Z\gamma$ for all neutral scalars, $W\gamma$ for all singly-charged)

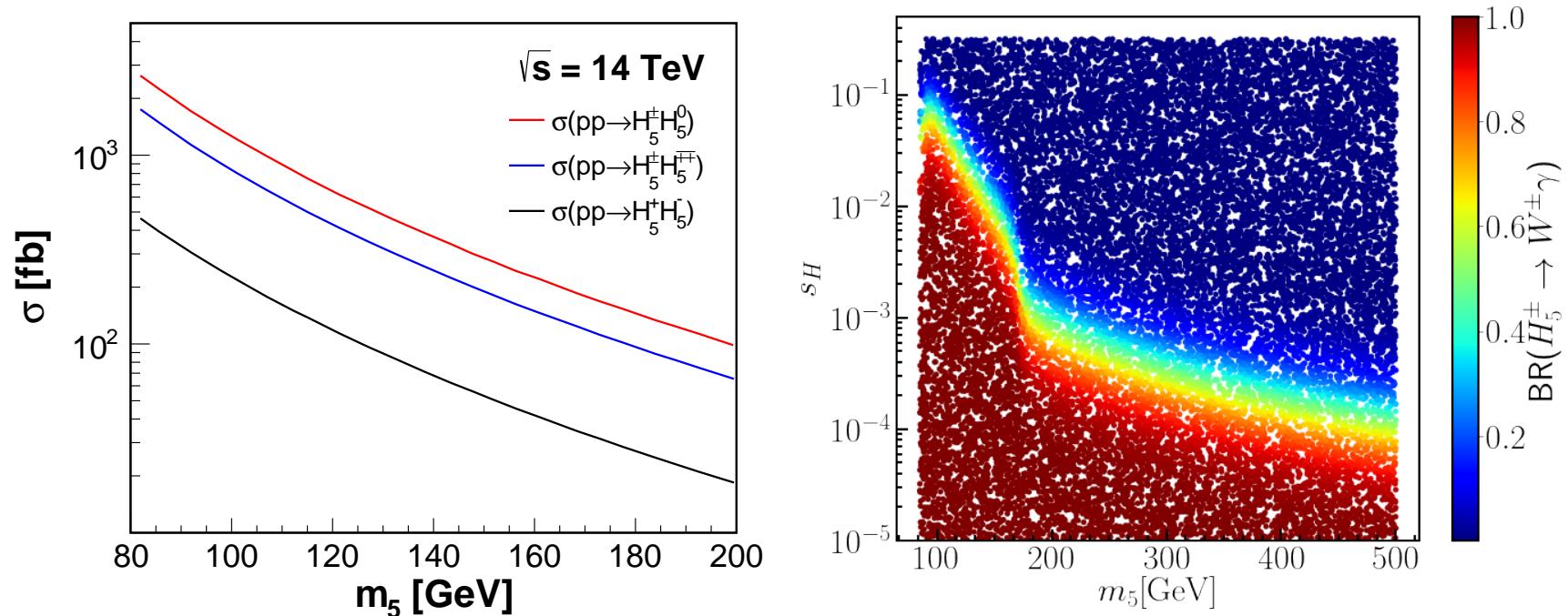
GMCALC 1.4.0: dedicated update to generate param_card.dat file including effective couplings for new EFT vertices
(also speed improvement for $H_i \rightarrow V_j^* V_k^*$, minor bug fixes)

Our simulation: MadGraph5 (LO) \rightarrow Pythia \rightarrow Delphes

H_5^\pm production and decays



Drell-Yan cross sections depend only on gauge couplings, m_5



LO QCD using MadGraph5, NNPDF23

$M_2 = 40$ GeV

Backgrounds:

$W^\pm\gamma$ (biggest, but killed well by cuts)

$t\bar{t}\gamma$ (biggest after cuts)

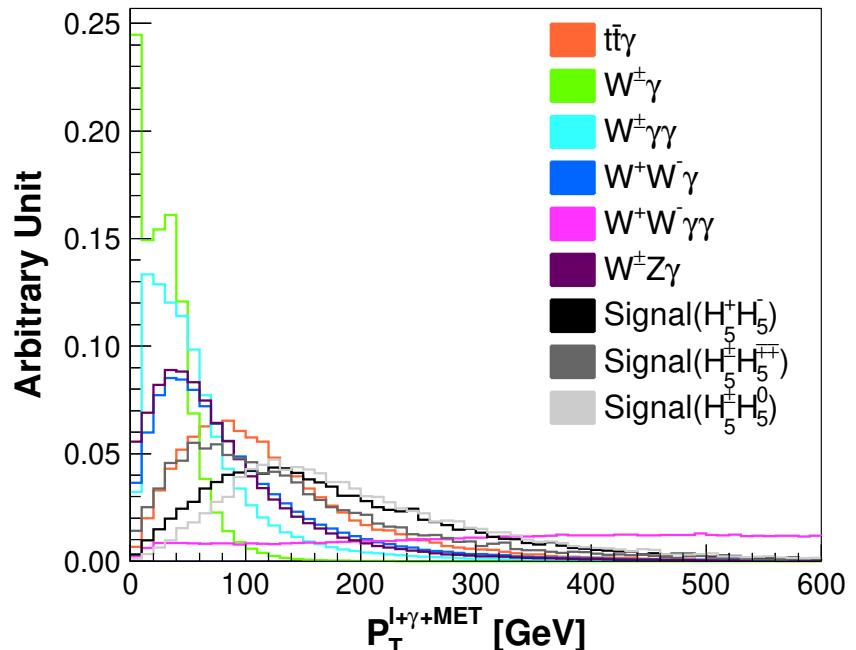
$W^+W^-\gamma$ (second biggest after cuts)

$W^+W^-\gamma\gamma$ (3rd biggest after cuts)

$W^\pm\gamma\gamma$

$W^\pm Z\gamma$

$$m_5 = 150 \text{ GeV} \rightarrow$$



Cuts:

≥ 1 lepton $p_T > 25 \text{ GeV}$, $|\eta| < 2.5$

≥ 1 photon $p_T > 25 \text{ GeV}$, $|\eta| < 2.5$

≤ 2 jets with $p_T > 20 \text{ GeV}$

Zero jets tagged as b -jets

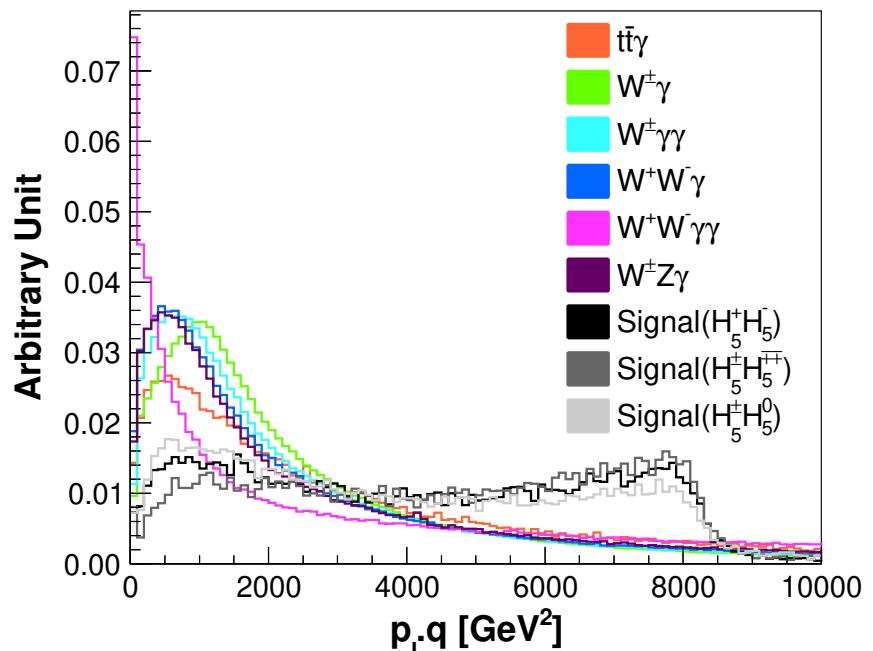
Optimize range for each mass:

MET

H_T (visible objects)

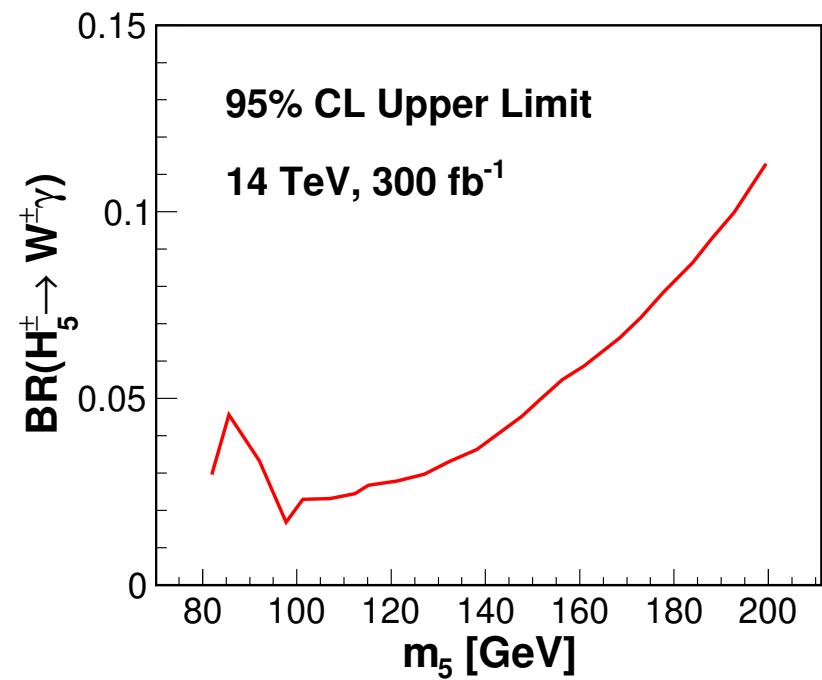
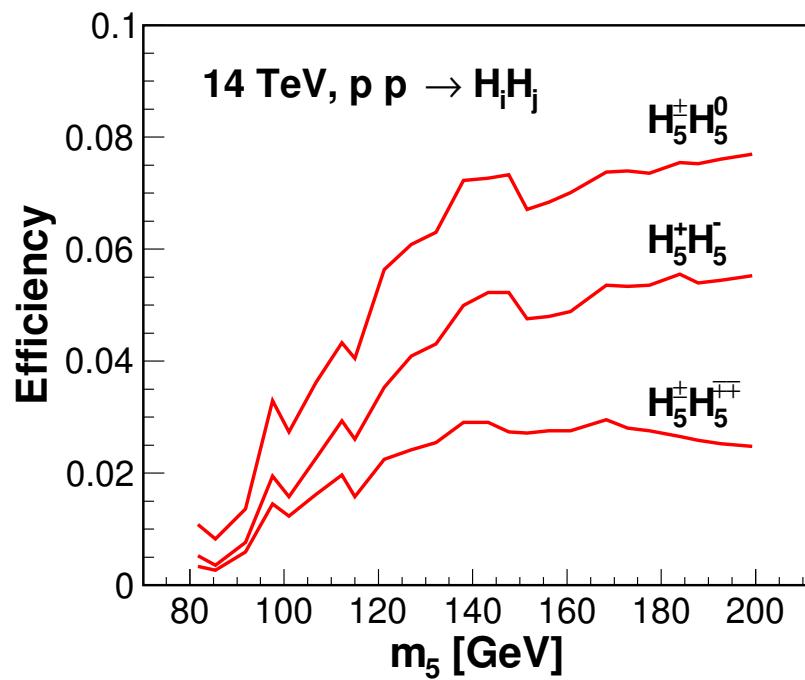
$p_T(\ell + \gamma + \text{MET})$

$p_\ell \cdot q$

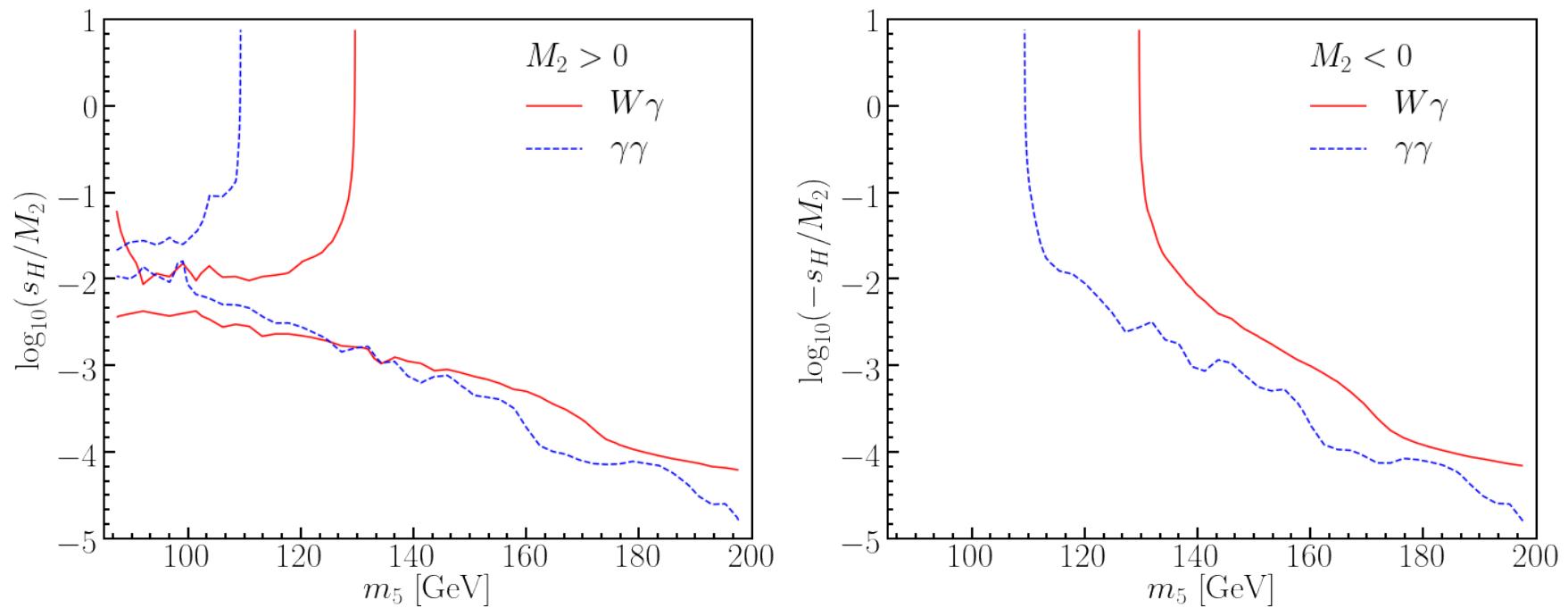


Results

$$\begin{aligned}
 (\sigma \times \text{BR})_{\text{Fiducial}} &\equiv \epsilon_{H_5^\pm H_5^0} \sigma(pp \rightarrow H_5^\pm H_5^0) \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) \\
 &+ \epsilon_{H_5^\pm H_5^{\mp\mp}} \sigma(pp \rightarrow H_5^\pm H_5^{\mp\mp}) \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) \\
 &+ \epsilon_{H_5^+ H_5^-} \sigma(pp \rightarrow H_5^+ H_5^-) [2\text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) - \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma)^2]
 \end{aligned}$$



Results



Red: Expected 95%CL exclusion with 300 fb^{-1} at 14 TeV LHC
Note cancellation between scalar and gauge diagrams at $s_H/M_2 \sim 10^{-2}/\text{GeV}$

Blue: Existing 95%CL exclusion from diphoton resonance search
ATLAS 1407.6583, 20.3 fb $^{-1}$ at 8 TeV
recast to constrain $pp \rightarrow H_5^0 H_5^\pm$ (LO), $H_5^0 \rightarrow \gamma\gamma$.

Conclusions

$H^\pm \rightarrow W^\pm \gamma$ is an interesting search channel for fermiophobic H^\pm with mass below WZ threshold

LHC prospects are good: use distinctive $p_\ell \cdot q$ distribution.

- Exclude $m_5 \lesssim 130$ GeV for almost all s_H values with 300 fb^{-1}
- Exclude m_5 up to 200 GeV and beyond for very small s_H

Georgi-Machacek model implementation with EFT vertices for loop-induced couplings now available as UFO file for MadGraph5 (LO QCD); param_card.dat generated by GMCALC 1.4.0

Strong competition from $H^0 \rightarrow \gamma\gamma$, but only when H^\pm and H^0 have the same mass (as is the case in GM model)

BACKUP SLIDES

Optimized cuts for $m_5 = 150$ GeV

$$72 \text{ GeV} \leq \text{MET} \leq 220 \text{ GeV}$$

$$260 \text{ GeV} < H_T < 620 \text{ GeV}$$

$$100 \text{ GeV} < p_T(\ell + \gamma + \text{MET}) < 420 \text{ GeV}$$

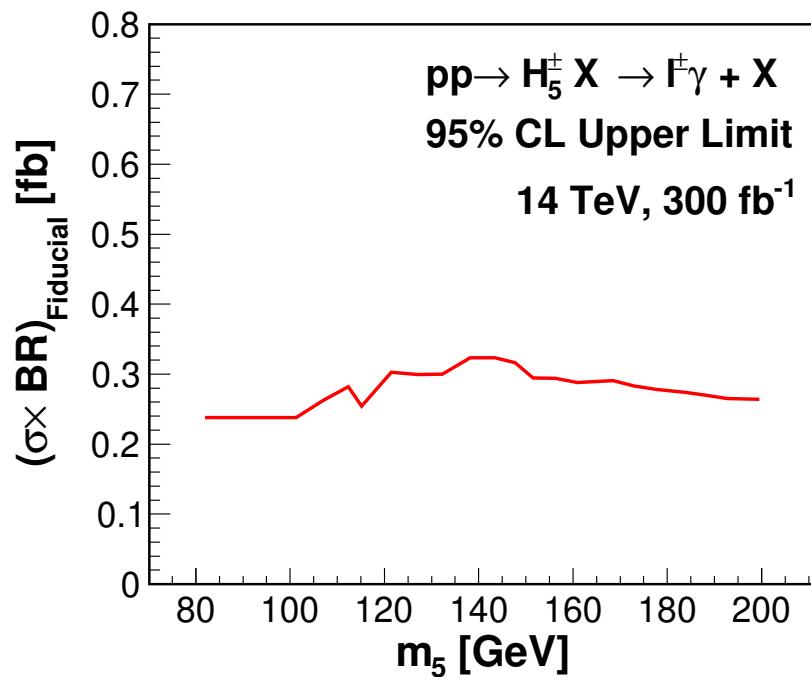
$$3300 \text{ GeV}^2 < p_\ell \cdot q < 8200 \text{ GeV}^2$$

Signal and background cross sections before and after cuts:

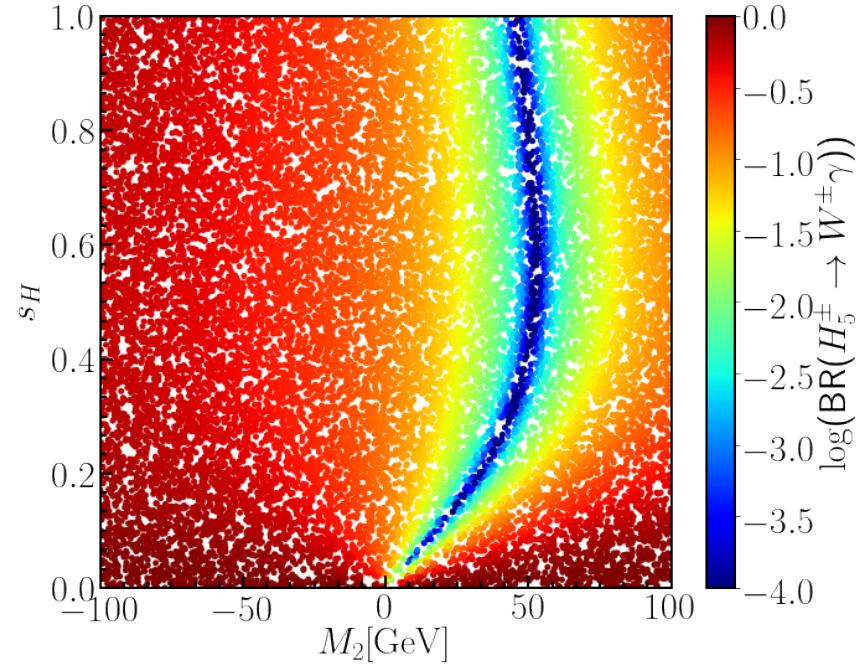
Process	$H_5^\pm H_5^0$	$H_5^\pm H_5^{\mp\mp}$	$H_5^+ H_5^-$	$t\bar{t}\gamma$	$W^\pm\gamma$	$W^\pm\gamma\gamma$	$W^+W^-\gamma$	$W^+W^-\gamma\gamma$	$W^\pm Z\gamma$
$\sigma \times \text{BR}$ [fb] (before cuts)	57.29	38.19	19.07	856	23000	30	120	65	25
$\epsilon \times \sigma \times \text{BR}$ [fb] (after cuts)	4.21	1.01	0.95	0.49	0.09	0.05	0.38	0.28	0.05

Expected upper limit on fiducial cross section

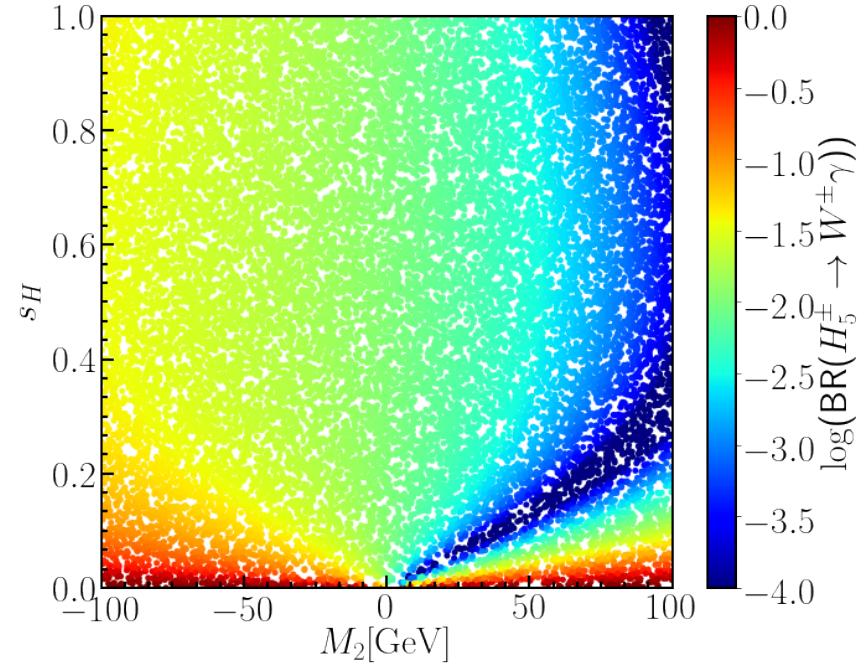
$$\begin{aligned}(\sigma \times \text{BR})_{\text{Fiducial}} &\equiv \epsilon_{H_5^\pm H_5^0} \sigma(pp \rightarrow H_5^\pm H_5^0) \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) \\&+ \epsilon_{H_5^\pm H_5^{\mp\mp}} \sigma(pp \rightarrow H_5^\pm H_5^{\mp\mp}) \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) \\&+ \epsilon_{H_5^+ H_5^-} \sigma(pp \rightarrow H_5^+ H_5^-) [2 \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma) - \text{BR}(H_5^\pm \rightarrow \ell^\pm \nu \gamma)^2]\end{aligned}$$



$H_5^+ \rightarrow W^+ \gamma$ branching ratio

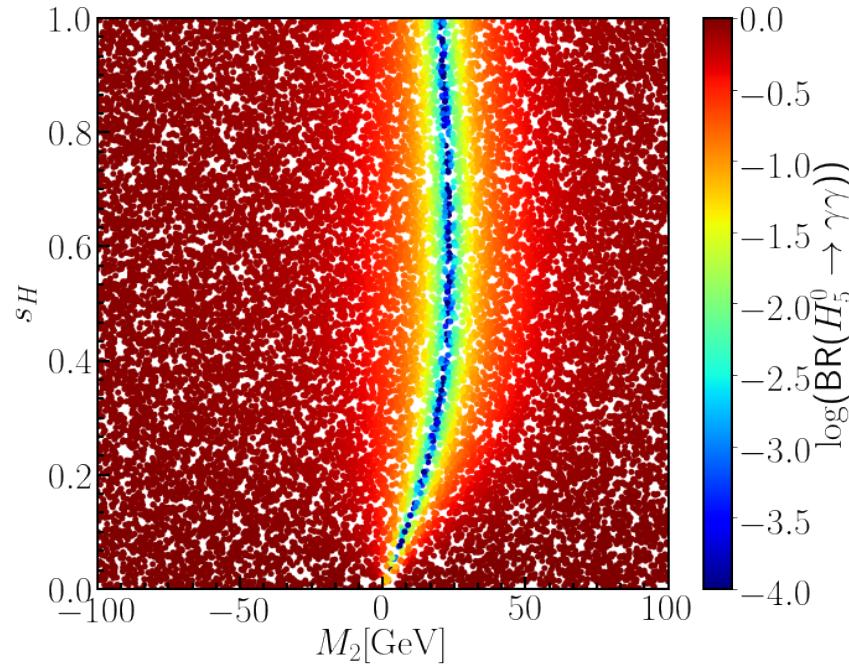


$m_5 = 100$ GeV

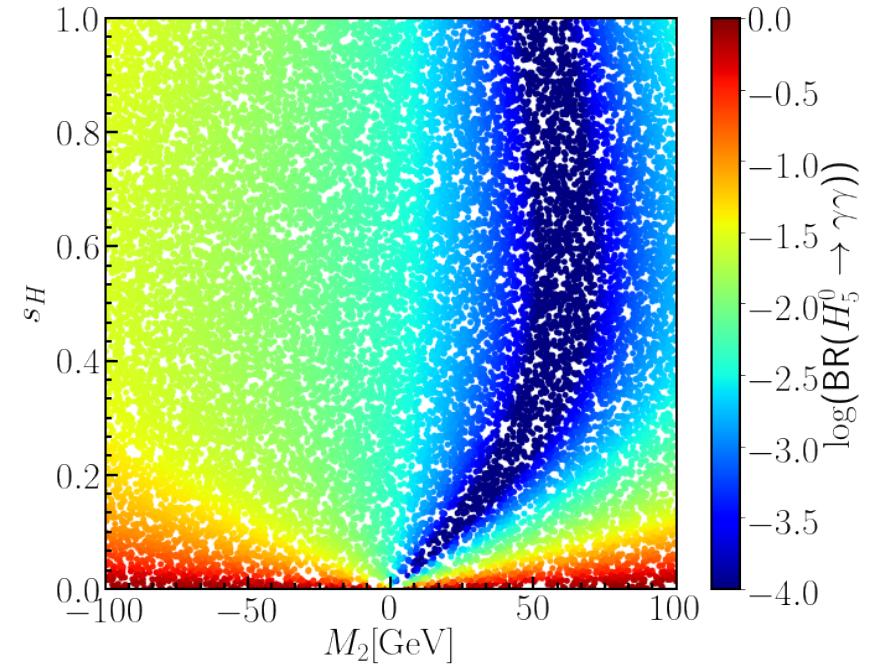


$m_5 = 150$ GeV

$H_5^0 \rightarrow \gamma\gamma$ branching ratio



$m_5 = 100$ GeV



$m_5 = 150$ GeV