## Formula sheet for final exam

This formula sheet, and copies of the front and back covers of Griffiths, will be provided with the exam.

Electric field due to a source charge q at the origin:

$$\vec{\mathbf{E}}(\vec{\mathbf{r}}) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$$
(1)

Force on test charge Q at  $\vec{\mathbf{r}}$  due to source charge q at the origin:

$$\vec{\mathbf{F}} = Q\vec{\mathbf{E}}(\vec{\mathbf{r}}) = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} \hat{\mathbf{r}}$$
(2)

Electric potential:

$$V(\vec{r}) = -\int_{\mathcal{O}}^{\vec{r}} \vec{\mathbf{E}} \cdot d\vec{\mathbf{l}}, \qquad \vec{\mathbf{E}} = -\vec{\nabla}V$$
(3)

Choosing  $\mathcal{O}$  at infinity, potential of a source charge q at the origin:

$$V(\vec{\mathbf{r}}) = \frac{1}{4\pi\epsilon_0} \frac{q}{r} \tag{4}$$

Gauss's law:

$$\oint_{S} \vec{\mathbf{E}} \cdot d\vec{\mathbf{a}} = \frac{Q_{\text{encl}}}{\epsilon_{0}}, \qquad \qquad \vec{\nabla} \cdot \vec{\mathbf{E}} = \frac{\rho}{\epsilon_{0}} \tag{5}$$

Poisson's equation (reduces to Laplace's equation for  $\rho = 0$ ):

$$\nabla^2 V = -\frac{\rho}{\epsilon_0} \tag{6}$$

Capacitance, definition of:

$$C = \frac{Q}{\Delta V}, \qquad C = \frac{A\epsilon_0}{d} \text{ for parallel plate}$$
(7)

Electrostatic energy for point charges and continuous charge distributions:

$$W = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \sum_{j>i}^n \frac{q_i q_j}{|\vec{r_i} - \vec{r_j}|}, \qquad \qquad W = \frac{1}{2} \int \rho \, V d\tau = \frac{\epsilon_0}{2} \int_{\text{all space}} E^2 d\tau \tag{8}$$

$$W = \frac{1}{2}C(\Delta V)^2 \text{ stored in a capacitor}$$
(9)

Discontinuities across a surface charge:

$$\vec{\mathbf{E}}_{\text{above}} - \vec{\mathbf{E}}_{\text{below}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}},\tag{10}$$

$$V_{\text{above}} = V_{\text{below}}, \qquad \qquad \frac{\partial V_{\text{above}}}{\partial n} - \frac{\partial V_{\text{below}}}{\partial n} = -\frac{\sigma}{\epsilon_0}$$
(11)

Electrostatic pressure (on surface of conductor):

$$\vec{f} = \frac{\sigma^2}{2\epsilon_0}\hat{n} = \frac{\epsilon_0}{2}E^2\hat{n} \tag{12}$$

Method of Images for a grounded conducting sphere with radius R and a point charge q a distance a from the centre of the sphere:

$$q' = -\frac{R}{a}q, \qquad b = \frac{R^2}{a} \tag{13}$$

where q' is a distance b from the centre of the sphere. For a conducting sphere at potential  $V_0$ , add a second image charge q'' at the centre of the sphere such that

$$V_0 = \frac{1}{4\pi\epsilon_0} \frac{q''}{R} \tag{14}$$

Solutions of Laplace's equation:

$$V(r,\theta) = \sum_{l=0}^{\infty} \left( A_l r^l + \frac{B_l}{r^{l+1}} \right) P_l(\cos\theta)$$
(15)

$$V(s,\phi) = A_0 \ln s + B_0 + \sum_{m=1}^{\infty} \left( A_m \cos m\phi + B_m \sin m\phi \right) \left( s^m + C_m s^{-m} \right)$$
(16)

Legendre polynomials:

$$P_{0}(x) = 1 \qquad P_{3}(x) = (5x^{3} - 3x)/2 P_{1}(x) = x \qquad P_{4}(x) = (35x^{4} - 30x^{2} + 3)/8 P_{2}(x) = (3x^{2} - 1)/2 \qquad P_{5}(x) = (63x^{5} - 70x^{3} + 15x)/8 e a dipole at the origin:$$

Potential due to a dipole at the origin:

$$V(\vec{\mathbf{r}}) = \frac{1}{4\pi\epsilon_0} \frac{\hat{\mathbf{r}} \cdot \vec{\mathbf{p}}}{r^2}$$
(17)

Bound charge in polarized materials:

$$\rho_b = -\vec{\nabla} \cdot \vec{\mathbf{P}} \qquad \sigma_b = \vec{\mathbf{P}} \cdot \hat{\mathbf{n}} \tag{18}$$

 $\vec{\mathbf{D}}$  field:

$$\vec{\mathbf{D}} \equiv \epsilon_0 \vec{\mathbf{E}} + \vec{\mathbf{P}} \qquad \vec{\nabla} \cdot \vec{\mathbf{D}} = \rho_f \qquad \oint \vec{\mathbf{D}} \cdot d\vec{\mathbf{a}} = Q_{f,\text{encl}} \tag{19}$$

Linear dielectrics: electric susceptibility  $\chi_e$ 

$$\vec{\mathbf{P}} = \epsilon_0 \chi_e \vec{\mathbf{E}} \qquad \vec{\mathbf{D}} = \epsilon_0 (1 + \chi_e) \vec{\mathbf{E}} \equiv \epsilon_0 \epsilon_r \vec{\mathbf{E}} \equiv \epsilon \vec{\mathbf{E}}$$
(20)

Energy

$$W = \frac{1}{2} \int \vec{\mathbf{D}} \cdot \vec{\mathbf{E}} \, d\tau \qquad \text{In vacuum, } W = \frac{\epsilon_0}{2} \int E^2 \, d\tau \qquad (21)$$

Lorentz force on point charge or current

$$\vec{\mathbf{F}} = q \, \vec{\mathbf{v}} \times \vec{\mathbf{B}} \qquad \vec{\mathbf{F}} = \int (\vec{\mathbf{J}} \times \vec{\mathbf{B}}) \, d\tau \qquad (22)$$

Biot-Savart law

$$\vec{\mathbf{B}}(\vec{\mathbf{r}}) = \frac{\mu_0}{4\pi} \int \frac{\vec{\mathbf{J}}(\vec{\mathbf{r}}') \times (\vec{\mathbf{r}} - \vec{\mathbf{r}}')}{|\vec{\mathbf{r}} - \vec{\mathbf{r}}'|^3} d\tau' \qquad \qquad \vec{\mathbf{B}}(\vec{\mathbf{r}}) = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{\mathbf{l}'} \times (\vec{\mathbf{r}} - \vec{\mathbf{r}'})}{|\vec{\mathbf{r}} - \vec{\mathbf{r}'}|^3}$$
(23)

Ampère's law (magnetostatics)

$$\vec{\nabla} \times \vec{\mathbf{B}} = \mu_0 \vec{\mathbf{J}} \qquad \oint \vec{\mathbf{B}} \cdot d\vec{\mathbf{l}} = \mu_0 I_{\text{encl}}$$
(24)

Vector potential

$$\vec{\mathbf{B}} = \vec{\nabla} \times \vec{\mathbf{A}} \qquad \text{If } \vec{\nabla} \cdot \vec{\mathbf{A}} = 0, \text{ then } \nabla^2 \vec{\mathbf{A}} = -\mu_0 \vec{\mathbf{J}} \text{ and } \vec{\mathbf{A}}(\vec{\mathbf{r}}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}')}{|\vec{\mathbf{r}} - \vec{\mathbf{r}'}|} d\tau' \qquad (25)$$

Boundary conditions

$$\vec{\mathbf{B}}_{\text{above}} - \vec{\mathbf{B}}_{\text{below}} = \mu_0(\vec{\mathbf{K}} \times \hat{\mathbf{n}})$$
(26)

 $\rightarrow \rightarrow$ 

Torque and force on a magnetic dipole

$$\vec{\tau} = \vec{\mathbf{m}} \times \vec{\mathbf{B}} \qquad \vec{\mathbf{F}} = \vec{\nabla} (\vec{\mathbf{m}} \cdot \vec{\mathbf{B}})$$
(27)

Vector potential due to a magnetic dipole at the origin

$$\vec{\mathbf{A}}(\vec{\mathbf{r}}) = \frac{\mu_0}{4\pi} \frac{\vec{\mathbf{m}} \times \hat{\mathbf{r}}}{r^2}$$
(28)

Bound current in magnetized materials

$$\vec{\mathbf{J}}_b = \vec{\nabla} \times \vec{\mathbf{M}} \qquad \quad \vec{\mathbf{K}}_b = \vec{\mathbf{M}} \times \hat{\mathbf{n}}$$
(29)

 $\vec{\mathbf{H}}$  field

$$\vec{\mathbf{H}} \equiv \frac{1}{\mu_0} \vec{\mathbf{B}} - \vec{\mathbf{M}} \qquad \vec{\nabla} \times \vec{\mathbf{H}} = \vec{\mathbf{J}}_f \qquad \oint \vec{\mathbf{H}} \cdot d\vec{\mathbf{l}} = I_{f,\text{encl}}$$
(30)

Linear magnetic media: magnetic susceptibility  $\chi_m$ 

$$\vec{\mathbf{M}} = \chi_m \vec{\mathbf{H}} \qquad \vec{\mathbf{B}} = \mu_0 (1 + \chi_m) \vec{\mathbf{H}} \equiv \mu \vec{\mathbf{H}}$$
(31)

Ohm's law and emf (here  $\sigma \equiv$  conductivity)

$$\vec{\mathbf{J}} = \sigma \vec{\mathbf{f}} = \sigma (\vec{\mathbf{E}} + \vec{\mathbf{v}} \times \vec{\mathbf{B}}) \qquad V = IR \qquad \mathcal{E} \equiv \oint \vec{\mathbf{f}} \cdot d\vec{\mathbf{l}}$$
(32)

Power

$$P = I\Delta V \tag{33}$$

Motional emf, induction, and Faraday's law

$$\mathcal{E} = -\frac{d\Phi_m}{dt} \qquad \vec{\nabla} \times \vec{\mathbf{E}} = -\frac{\partial \vec{\mathbf{B}}}{\partial t} \qquad \oint \vec{\mathbf{E}} \cdot d\vec{\mathbf{l}} = -\int \frac{\partial \vec{\mathbf{B}}}{\partial t} \cdot d\vec{\mathbf{a}}$$
(34)

Lenz's law: "Nature abhors a change in flux." Mutual inductance and self-inductance

$$\Phi_2 = MI_1 \qquad \Phi_1 = MI_2 \qquad \Phi = LI \qquad \mathcal{E} = -L\frac{dI}{dt} \tag{35}$$

Energy

$$W = \frac{1}{2\mu_0} \int B^2 d\tau \tag{36}$$

Continuity equation and Maxwell's fix to Ampère's law

$$\vec{\nabla} \cdot \vec{\mathbf{J}} = -\frac{\partial \rho}{\partial t} \qquad \qquad \vec{\nabla} \times \vec{\mathbf{B}} = \mu_0 \vec{\mathbf{J}} + \mu_0 \epsilon_0 \frac{\partial \vec{\mathbf{E}}}{\partial t} \tag{37}$$